

Wakefield Effects in the Spreader of LCLS-II LCLS-II TN-15-04

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Introduction

• In LCLS-II the trajectories of the hard and soft x-ray beam will be separated by a kicker, with one type of bunch kicked vertically in one direction, and the other type kicked in the opposite direction.

• With only one kicker chamber, both beams will on average be offset from the axis of the (round) beam pipe chamber. This will induce transverse (dipole) wakefields that generate a transverse kick that varies along the bunch, thus increasing the projected emittance. Is the problem severe enough that two separate chambers are needed?

• The chamber walls are ceramic that is copper coated, and the resistive wall wake will be excited. The chamber is round with radius a = 5 mm, length L = 8 m. The beam enter on axis, follows a vertical parabolic trajectory, and exits offset by 3 mm from the axis. The average offset is $\langle y \rangle = -1$ mm

 \bullet Here I will calculate the transverse wake effect, as well as that of the longitudinal wake

Introduction Cont'd

• In LCLS-II, after BC2, the beam (longitudinal) distribution is approximately uniform, with peak current I = 1 kA. The nominal bunch charge Q = 100 pC, the largest charge is 300 pC. The total bunch length for the two cases is 30, 90 μ m

• Calculate the full rw wake, ac model with conductivity $\sigma_c = 6.0 \times 10^7 \ \Omega^{-1} m^{-1}$, relaxation time $\tau = 2.7 \times 10^{-14}$ s (Cu). [See e.g. K. Bane and M. Sands, AIP Conf. Proc. 367 (1996) 131]

• Currently in the optics β_y varies from 112 m to 237 m; β_y will be reduced to 33 m–Yuri N is working on it

• Another effect that needs to be investigated (–Tor) for the kicker (not a wake effect, and it will not be done in this talk), is the stay clear. When the beam leaves the chamber it passes at a distance of 2 mm from the wall. With the optics assumed here, the beam size $\sigma_{\gamma} \sim 50 \ \mu$ m, that distance is equivalent to $40\sigma_{\gamma}$

Resistive Wall Wakes

• The rw wake scale factor

$$s_0 = \left(\frac{2a^2}{Z_0\sigma_c}\right)^{1/3}$$

with a the beam pipe radius, σ_c the dc conductivity, and $Z_0=377~\Omega.$ For kicker chamber $s_0=13~\mu\text{m}$

 \bullet For ac rw model, replace σ_c by $\tilde{\sigma}_c=\sigma_c/(1-\textit{ikc}\tau),$ with τ the relaxation time

• Longitudinal rw impedance as function of normalized wave number $\kappa = ks_0$:

$$Z_{I}(\kappa) = 2\left(\frac{s_{0}}{ca^{2}}\right)\left(\frac{2}{1-i}\frac{1}{\sqrt{\kappa}} - i\frac{\kappa}{2}\right)^{-1}$$

• Transverse (dipole) rw impedance

$$Z_{yd}(\kappa) = 4\left(\frac{s_0^2}{ca^4\kappa}\right)\left(\frac{2}{1-i}\frac{1}{\sqrt{\kappa}} - i\frac{\kappa}{2}\right)^{-1}$$

Resistive Wall Wakes Cont'd

• Longitudinal wake

$$W_l(s) = rac{c}{2\pi} \int_{-\infty}^{\infty} dk \, Z_l(k) e^{-iks}$$

Units are $[V/(pC \cdot m)]$

• For uniform bunch distribution with peak current *I*, in pipe of length *L*, voltage within bunch is given by

$$E_w(s) = -\frac{lL}{c} \int_0^s ds' W_l(s')$$

• The transverse wake is obtained similarly from Z_{yd} . The kick angle, for a beam of energy E, offset by y_0 , at the end of a pipe of length L is

$$\Delta y'(s) = \frac{elLy_0}{cE} \int_0^s ds' \, W_{yd}(s')$$

Longitudinal Effect



Figure: Rw wake induced energy change in LCLS-II spreader kicker chamber for a bunch with a uniform distribution, with peak current I = 1 kA. The bunch head is located at s = 0

At E = 4 GeV, the total spread in induced relative energy variation is $\delta E_w = 7 (9) \times 10^{-5}$, for Q = 100 (300) pC

Transverse Wake



Figure: Transverse rw (point charge) wake in LCLS-II spreader kicker chamber. Slope at origin is $W'_{v}(s) = 2Z_0c/\pi a^4$

Transverse Effect

• Calculate vertical kick for a bunch passing through the pipe offset by y = 1 mm: Here E = 4 GeV, $\beta_y = 33$ m, $\epsilon_n = 0.5$ µm



Figure: Rw wake induced transverse (dipole) kick in LCLS-II spreader kicker chamber for a bunch with a uniform distribution with peak current I = 1 kA that is moving on the design orbit. Plotted is $\Delta y'/\sigma_{y'}$

• Due to the off-axis orbit, the last slice in the 100 pC or 300 pC bunches will be displaced by ~ 10% of $\sigma_{y'}$ at the end of the kicker, assuming E=4 GeV, $\beta_y=33$ m, ε_n =0.5 μm . For larger β_y the effect will grow by $\beta_y^{1/2}$

• Paul E. pointed out that this a small effect in terms of the emittance. The effective emittance growth for the last slice is $\Delta \epsilon_y/\epsilon_y = \frac{1}{2} (\Delta y'/\sigma_{y'})^2 = 0.5\%$

• The longitudinal effect is a wake-induced total energy variation, for both the 100 pC and 300 pC cases, of $\Delta E_w/E \sim 10^{-4}$. The chirp is very non-linear and probably cannot be corrected