

Interstage Optics Design

for a PWFA Linear Collider

FACETII Workshop 2015 – Oct 14, 2015

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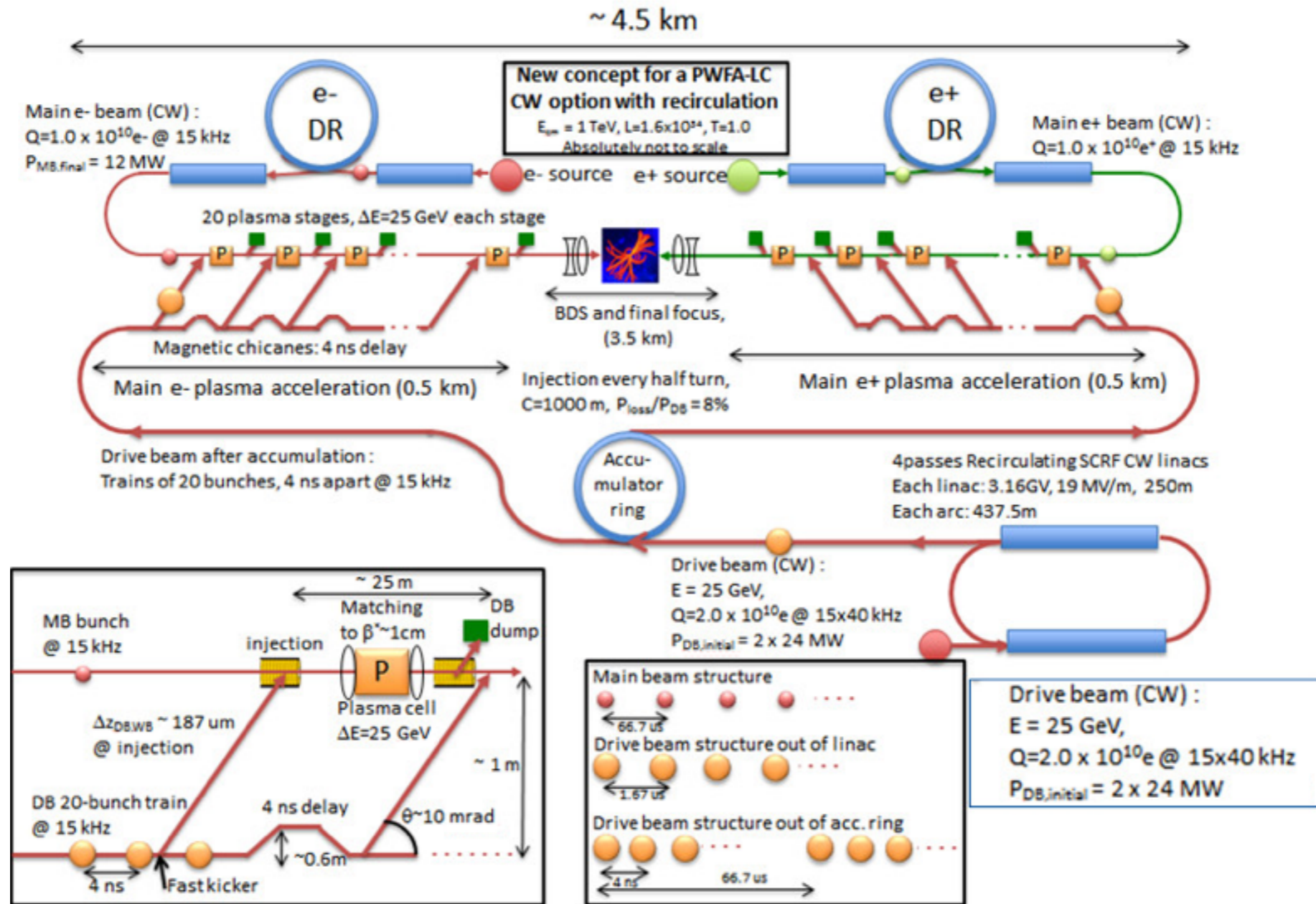
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Proposed layout of a PWFA Linear Collider*

- Basis for this study + source of parameters.

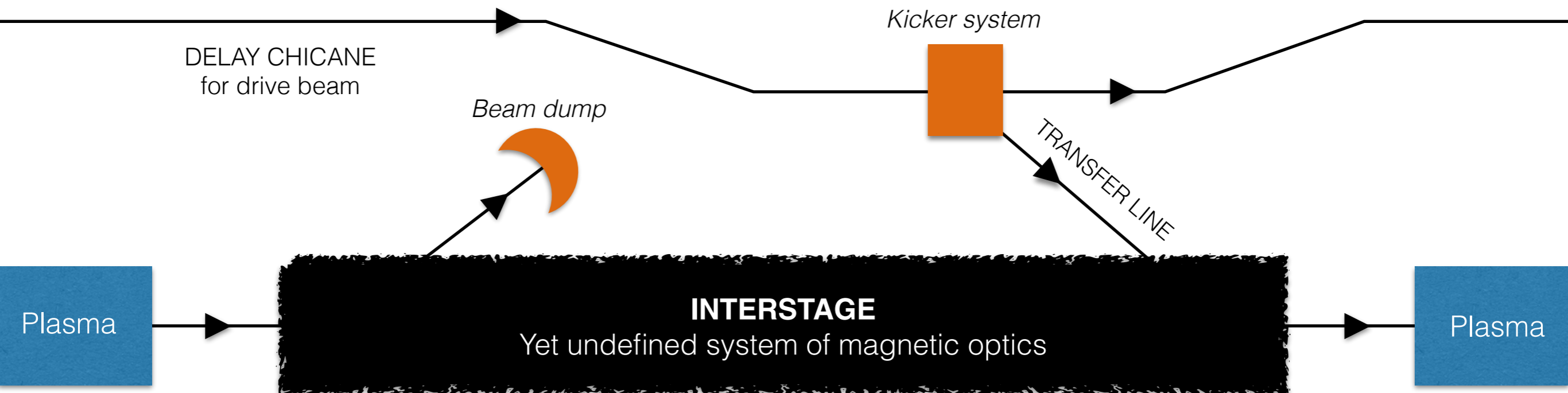


* E Adli, JP Delahaye, et al. (presented at IPAC'13)

Interstage optics

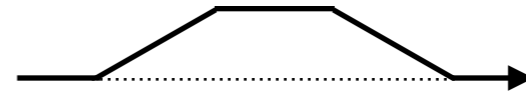
- Problem at hand:

Switch old with fresh drive beam,
keep the main beam focused
and preserve its emittance.



Formal requirements

- Drive beam injection/extraction
- Collinearity
- Beta function matching
- Dispersion cancellation
- Limit bunch lengthening (R_{56})
- Emittance preservation
 - limit chromaticity
 - limit synchrotron radiation
- Minimize length subject to all the above



$$\beta_x(L) = \beta_y(L) = \beta_{mat}$$

$$\alpha_x(L) = \alpha_y(L) = 0$$

$$D_x(L) = D'_x(L) = 0$$

$$R_{56}(L) \ll \frac{\sigma_z}{\delta} \approx 1 \text{ mm}$$

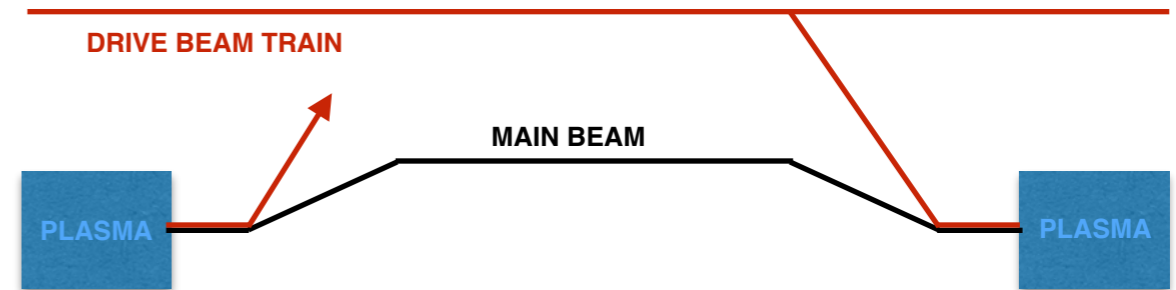
$$\frac{\Delta\epsilon}{\epsilon}(L) \ll 1\%$$

Injection/extraction of drive beams

- Using dipoles to create dispersion:
Separate beams spatially by energy.
- Injection and extraction are symmetric processes \Rightarrow mirror symmetric lattice.
- Injection/extraction dipoles must be first and last magnets, as main beam quads would destroy the drive beam.
- Important: Dipoles do not scale with main beam energy (only drive beam energy).
- Defines regimes:
 - Low energy ($E_{\text{main}} \approx E_{\text{drive}}$) \Rightarrow Dipoles “visible”
 - High energy ($E_{\text{main}} \gg E_{\text{drive}}$) \Rightarrow Dipoles “invisible”
- Scalings:
 - Dipole field & length: $L_{\text{dipole}}, B \sim \text{constant}$
 - SR power: $P_{\text{SR}} \sim E_{\text{main}}^2$
 - Main beam dispersion: $D_x \sim L_{\text{interstage}}/E_{\text{main}} \sim 1/\sqrt{E_{\text{main}}}$
(assuming focusing $\Rightarrow L_{\text{interstage}} \sim \sqrt{E_{\text{main}}}$)

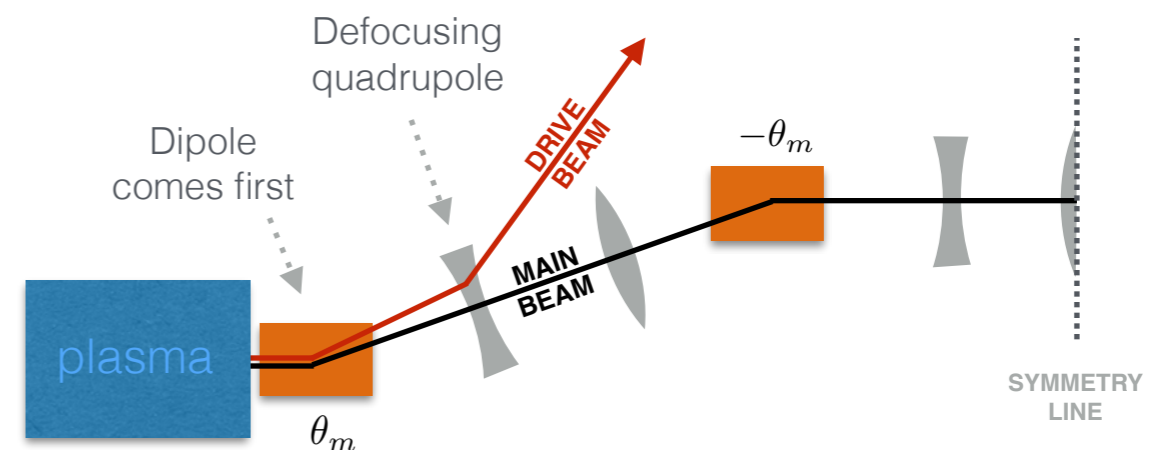
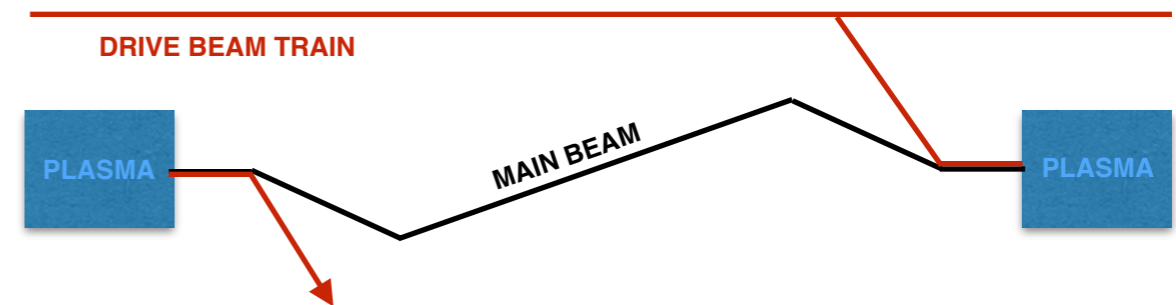
Option 1: “C-chicane”

Weaker bending



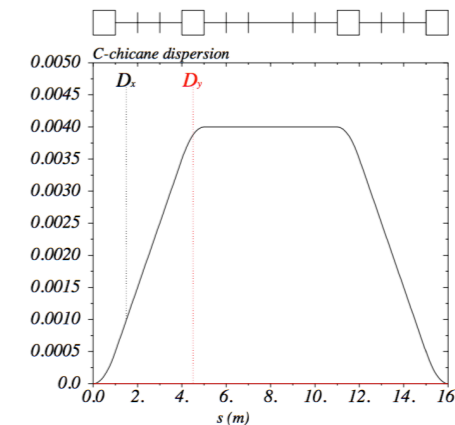
Option 2: “S-chicane”

Stronger bending, more space for beam dump.

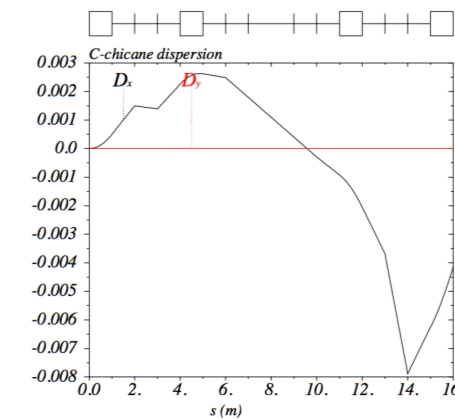


Dispersion cancellation

- Both S- and C-chicane has inherent dispersion cancellation. Quadrupoles change this.
- Cancel by:
 - Matching quadrupoles (not independent of beam matching)
 - Inserting extra dipoles (independent of beam matching)
- Low energy regime (large main beam dispersion) may require second order dispersion cancellation.
- Becomes less important with higher energies: **Scaling:** $D_x(s) \sim 1/\sqrt{E_{\text{main}}}$



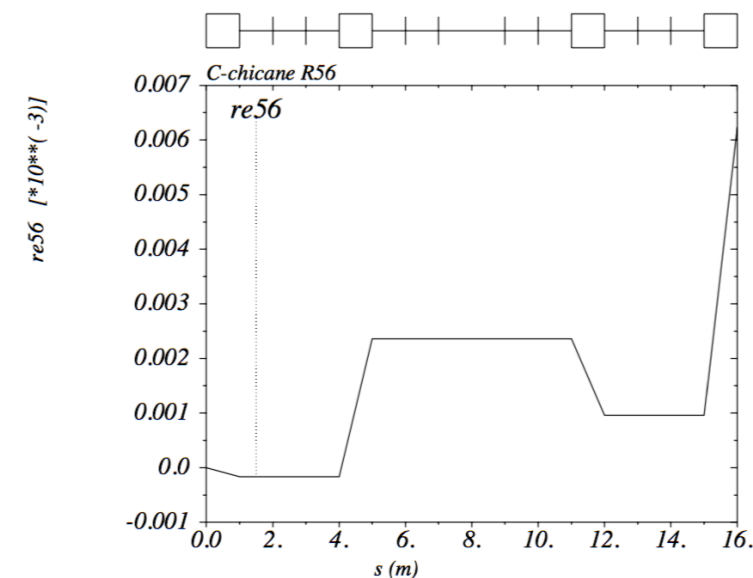
D_x
without
quads



D_x
with
quads

Limiting bunch lengthening (R_{56})

- Not big problem due to relatively weak dipoles.
- Becomes easier with higher energy. **Scaling:** $R_{56}(s) \sim D_x(s) \sim 1/\sqrt{E_{\text{main}}}$
- If necessary, a method for matching $R_{56} \approx 0$:
 - Low dispersion in dipoles.



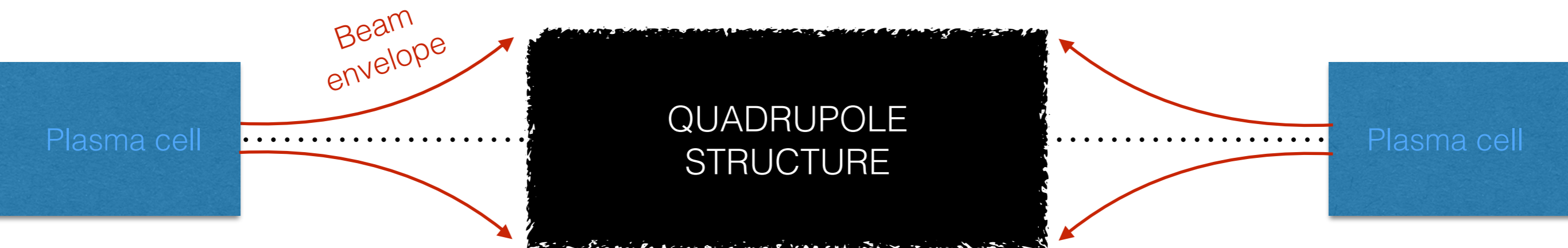
R₅₆
with
quads

$$R_{56}(L) \ll \frac{\sigma_z}{\delta} \approx 1 \text{ mm}$$

Matching beam to plasma

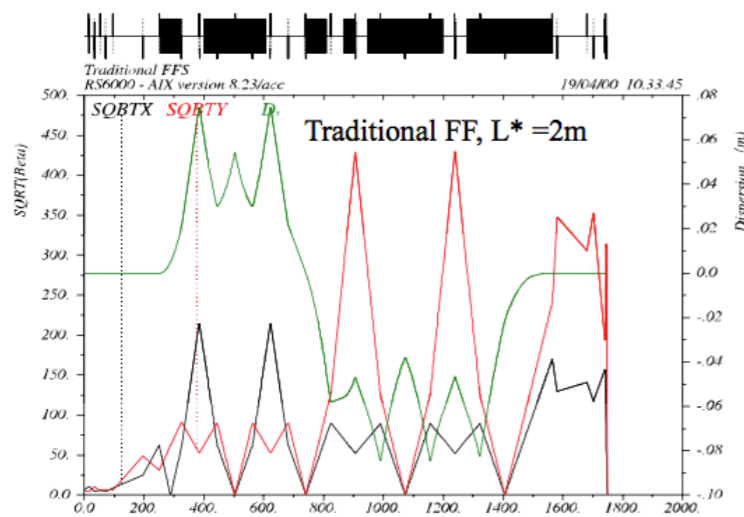
$$\alpha_x(L) = \alpha_y(L) = 0 \quad \beta_x(L) = \beta_y(L) = \beta_{mat}$$

- Matched beta in a plasma: $\beta_{matched} = \frac{\sqrt{2\gamma}}{k_p} = \sqrt{\frac{2\epsilon_0 E}{n_p e^2}} \Rightarrow 3.3 \text{ cm @ } E = 100 \text{ GeV, } n_p = 10^{16} \text{ cm}^{-3}$
- Naive matching is simple: **place quadrupoles, match strengths/separations.**
- Strong plasma channel focusing** \Rightarrow small plasma betas
 - \Rightarrow strongly diverging/large beams
 - \Rightarrow long/strong quadrupoles
 - \Rightarrow **large chromaticity (big challenge)**
- Requirements similar to those of a **final focus system, for every stage.**

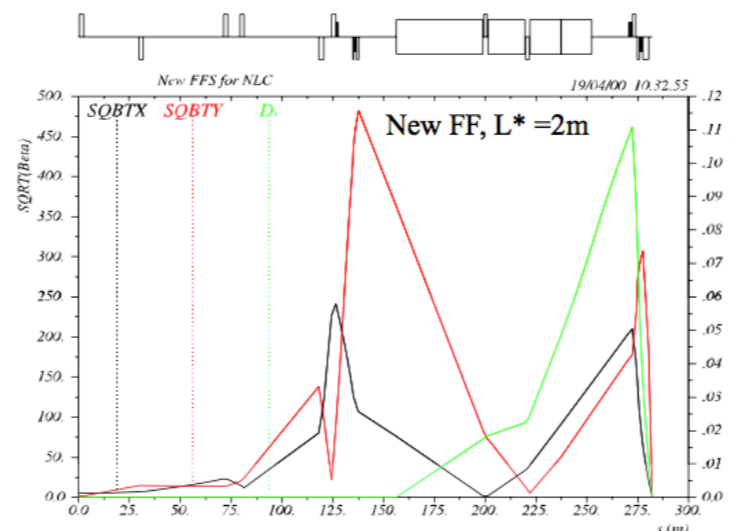


Chromaticity cancellation

- Small beams in plasma lead to large chromaticity.
- PWFALC study assumes a $\sim 1\%$ rms. energy spread.
- Conventionally corrected using sextupoles.
Seen in final focus systems (SLC, ILC...): very long lattices.
- We will consider a sextupole solution, and a novel solution.



Traditional FF @ 500GeV ~ 1750 m

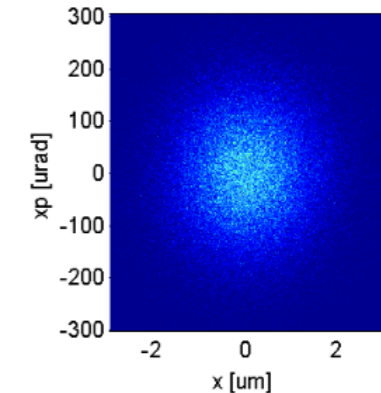


Improved FF @ 500GeV ~ 300 m

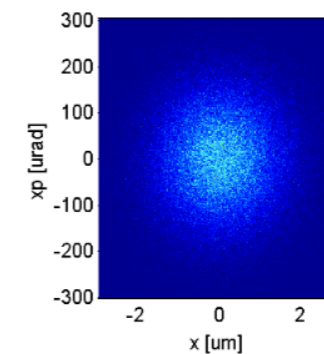
*Figures from "Beam Delivery & beam-beam" by **Andrei Seryi (SLAC)**

Example:

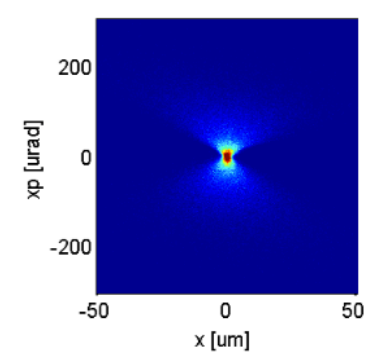
Initial phase space (x, x')



Final phase space
(no energy spread)



Final phase space
(1% energy spread)

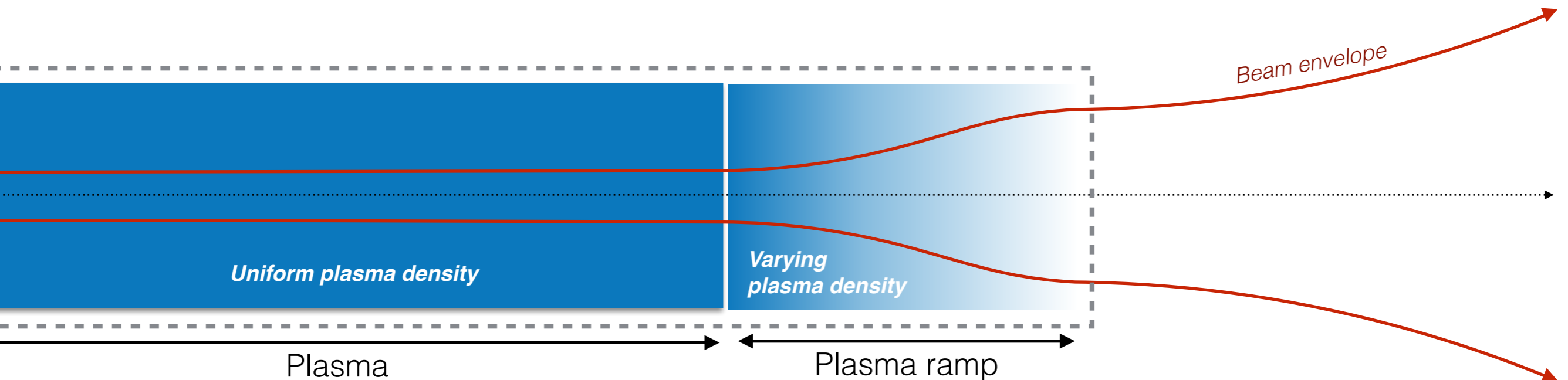


No chromaticity correction:
Emittance increases by many orders of magnitude

Plasma density ramp

* Plasma ramp paper by Xu et al. (2015): <http://arxiv.org/pdf/1411.4386v2.pdf>

- A plasma density ramp **mitigates the chromaticity problem.**
- Plasma ramp \Rightarrow larger/less diverging beam
 \Rightarrow smaller beam in quadrupoles \Rightarrow less chromaticity
- Any adiabatic density profile will do. Use pressure gradients or partial laser-ionization.
- **Scaling** (for magnification factor Π_p): $L_{\text{ramp}} \sim \Pi_p \sqrt{E_{\text{main}}}$



Conventional solution: Sextupoles

- Sextupole effect is stronger with:
 - Larger dispersion
 - Larger beam size
- Long lattices for geometric term cancellation.
- Introduces new problems:
 - Dipoles must ramp with main beam energy
 ⇒ Dispersion / R_{56} / SR scales poorly with energy
 - -I transforms require repeated sections
 ⇒ “unnecessarily” long lattices
 - Thick sextupoles (imperfect -I transforms)
 ⇒ geometric errors (emittance growth)
 - Sextupoles need large beam sizes
 ⇒ increased energy spread from SR (Oide effect)

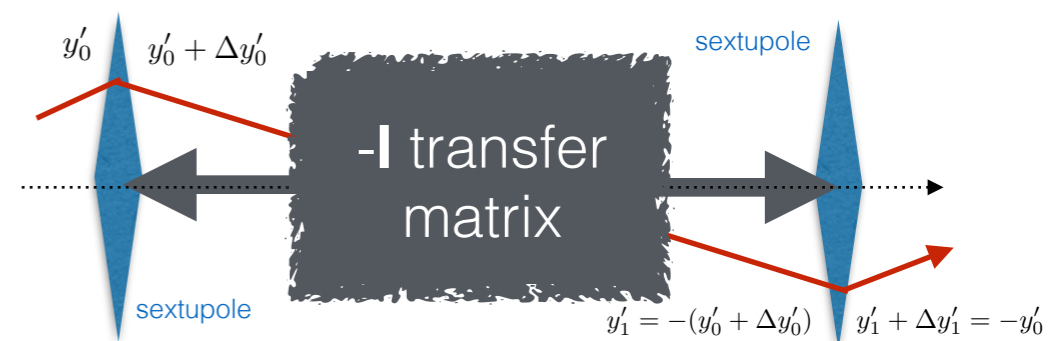
Sextupole B-fields:

Non-linear geometric terms **NEEDS TO BE SMALL** Non-linear chromatic term **NEED TO BE CANCELLED**

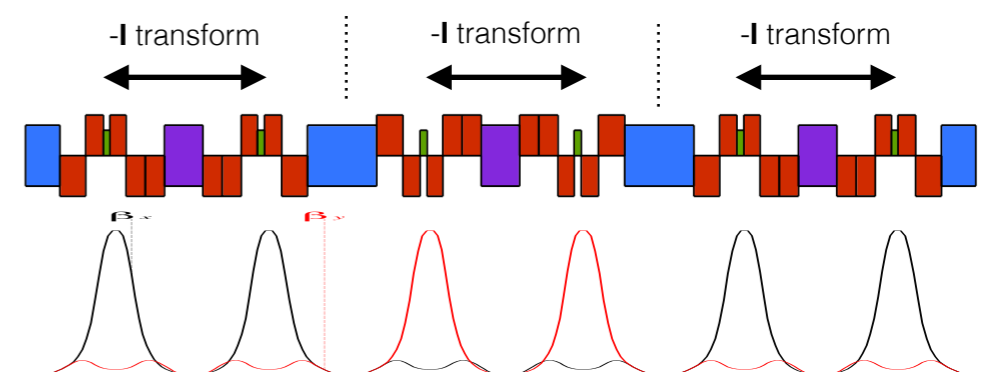
$$B_x \sim xy + \delta D_x y \quad B_y \sim \frac{1}{2}(x^2 - y^2) + x\delta D_x + \frac{1}{2}\delta^2 D_x^2$$

Linear chromatic terms **CORRECT CHROMATICITY**

Geometric term cancellation:



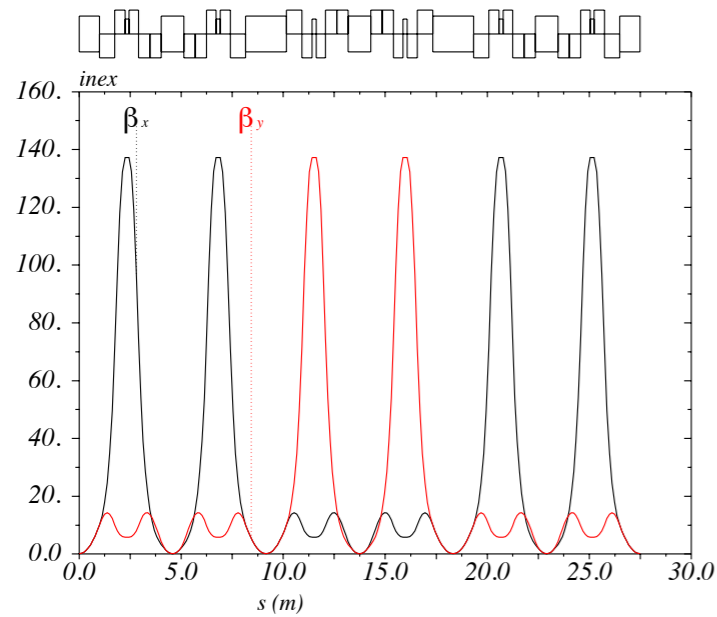
“Working” interstage using sextupoles:



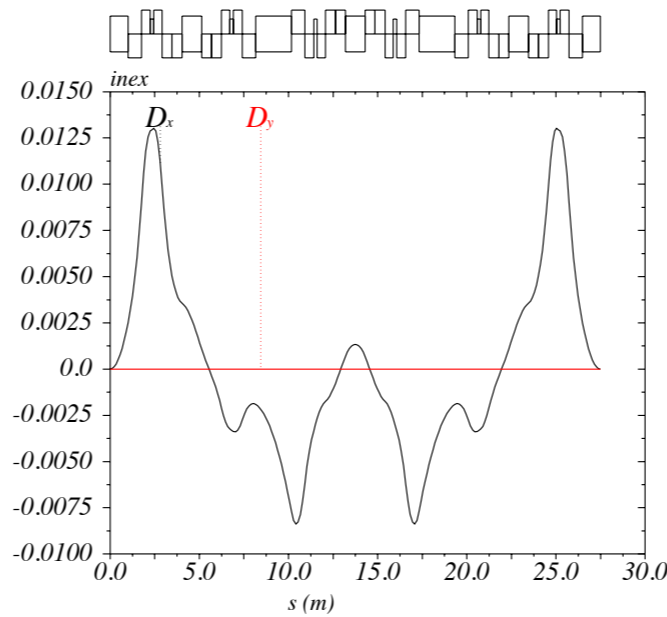
“Working” interstage using sextupoles

100 GeV, $\beta_{mat} = 0.1m$

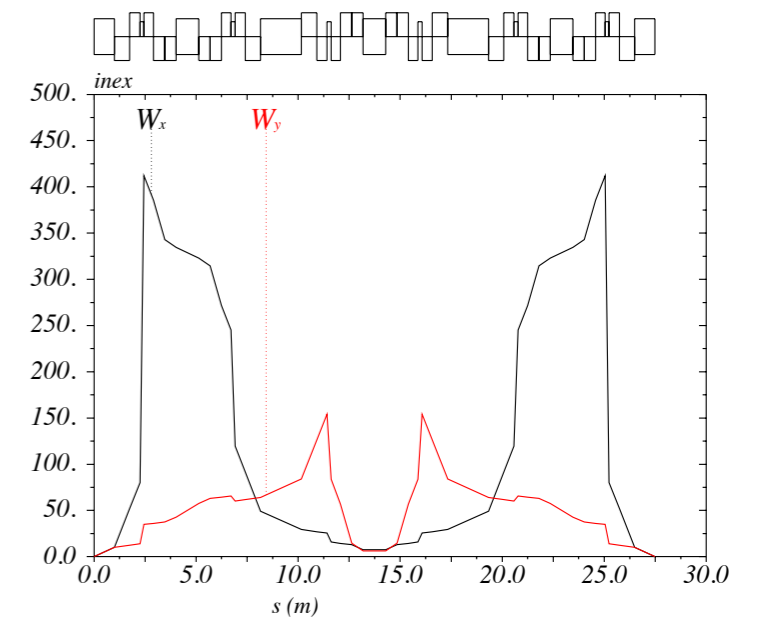
Beta functions



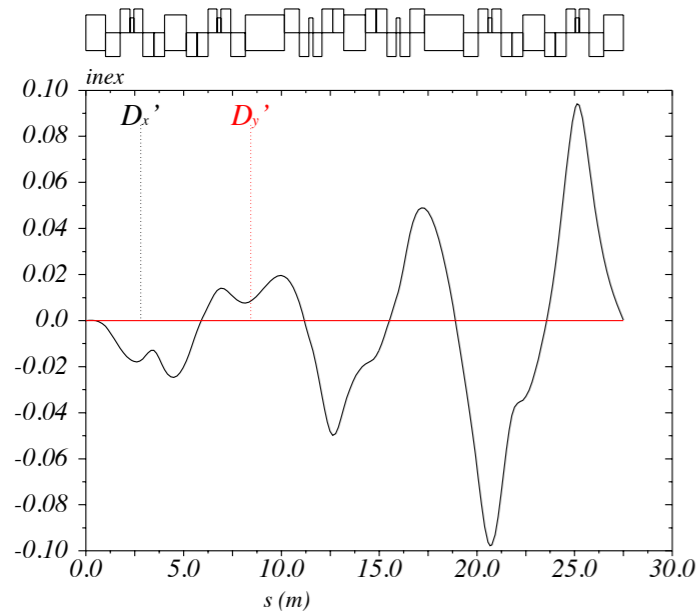
Dispersion



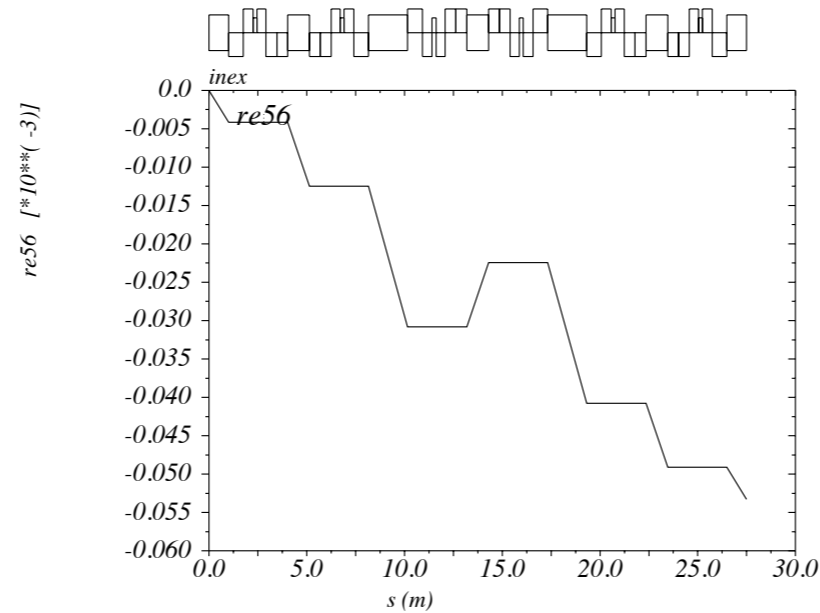
W-functions (chrom.)



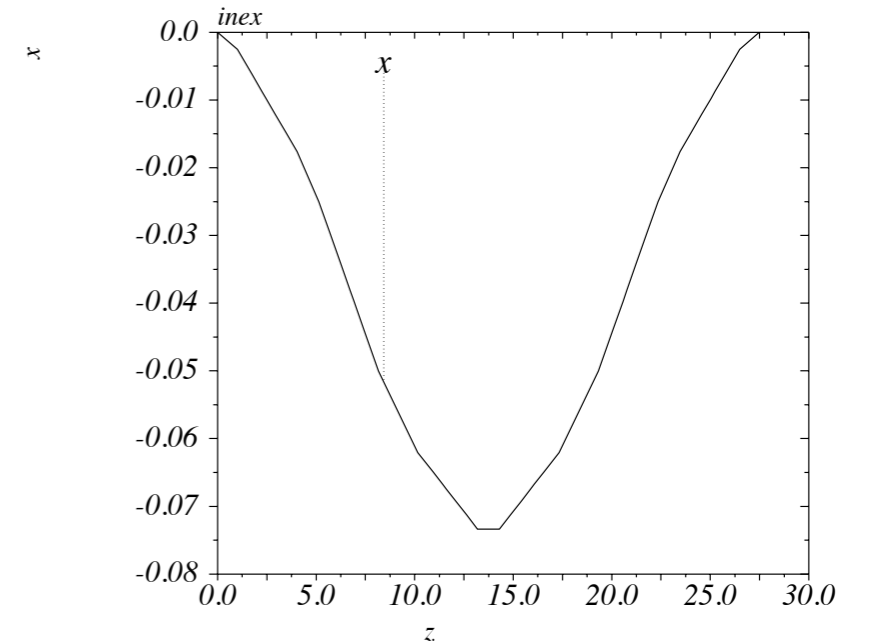
Second order dispersion



R_{56}



Footprint (x-z)



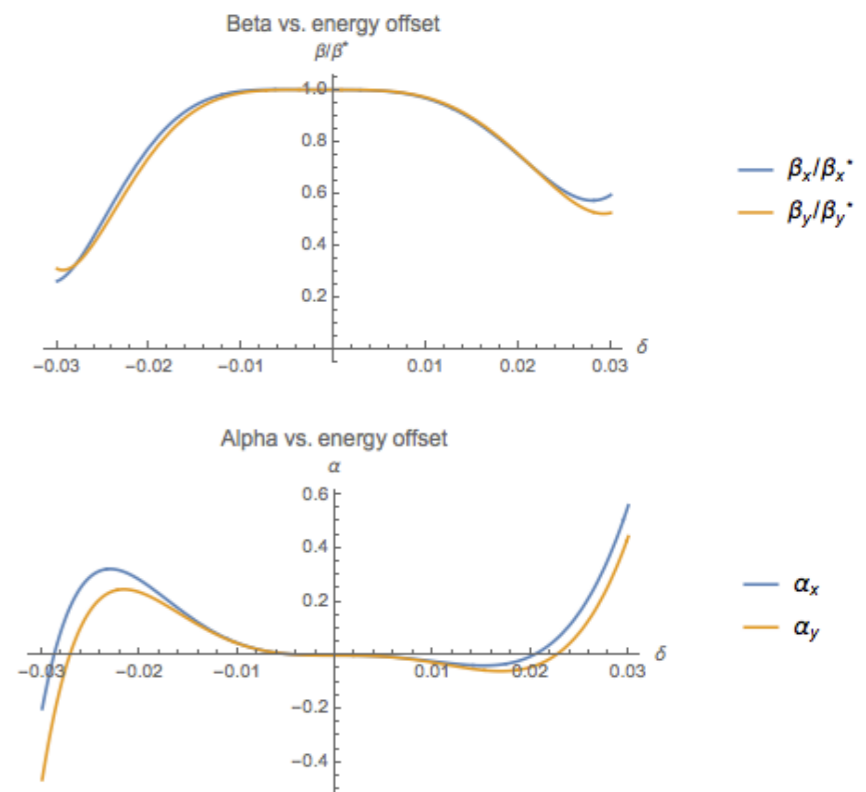
- This solution requires stronger sextupoles than currently manufacturable. Not a conceptual show-stopper.

Novel solution: Getting rid of sextupoles

- We are developing a new method for finding quadrupole-only lattices which cancel chromaticity to the required order in energy offset.
- Benefits of using quadrupoles only:
 - no geometric terms \Rightarrow keeps it linear
 - no -I transforms \Rightarrow much shorter
 - no ramping dipoles \Rightarrow better SR scaling
- Similar solutions in light optics: **Superachromats** (however, beam optics is x/y-asymmetric).
- Achieved by tailoring the energy offset (δ) expansion of β and α to be flat around $\delta = 0$.



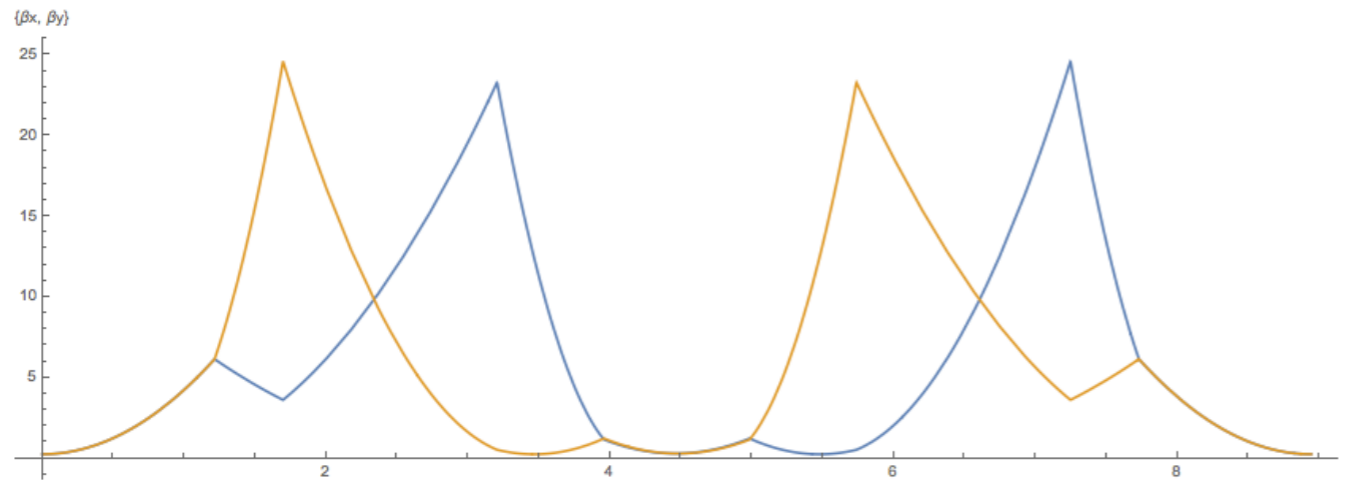
Simpler version used in light optics:
"Superachromat" by Carl Zeiss



β and α vs. energy offset δ .
Flat regions \Rightarrow no chromaticity

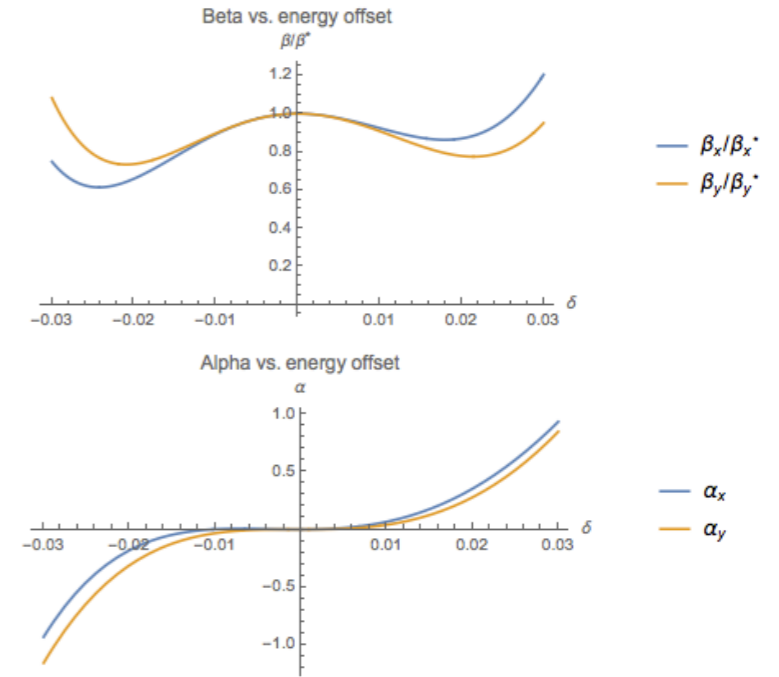
Examples of chromaticity-free quadrupole lattices

- 8 quads: cancel chromaticity to 1st order.

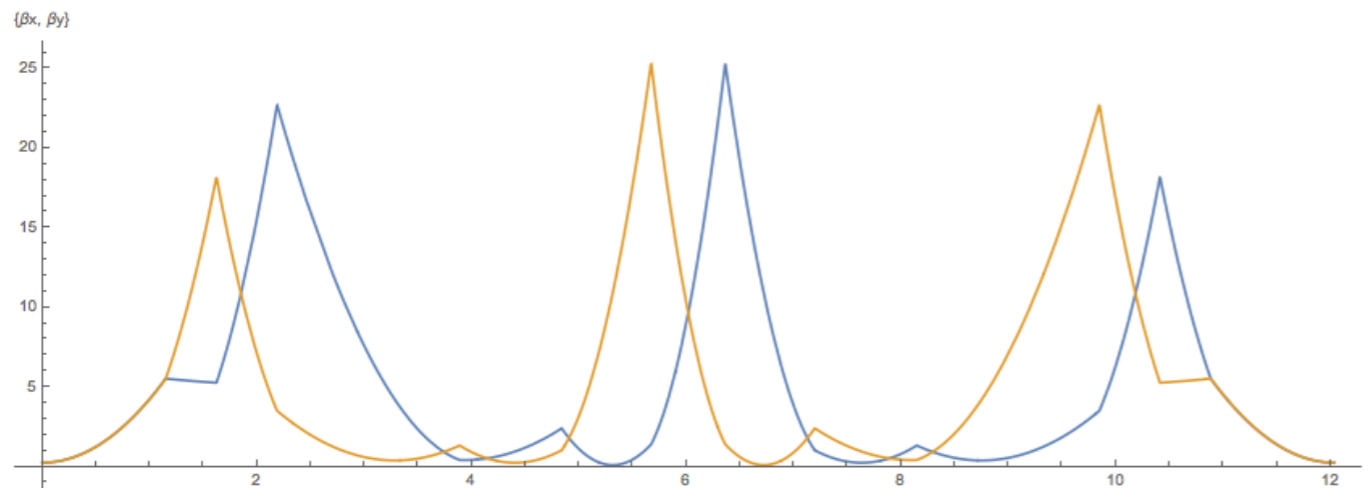


For 0.5% energy spread: $(\Delta\epsilon_x/\epsilon_x, \Delta\epsilon_y/\epsilon_y) = (0.000781, 0.000782)$

For 1.% energy spread: $(\Delta\epsilon_x/\epsilon_x, \Delta\epsilon_y/\epsilon_y) = (0.0141, 0.0142)$

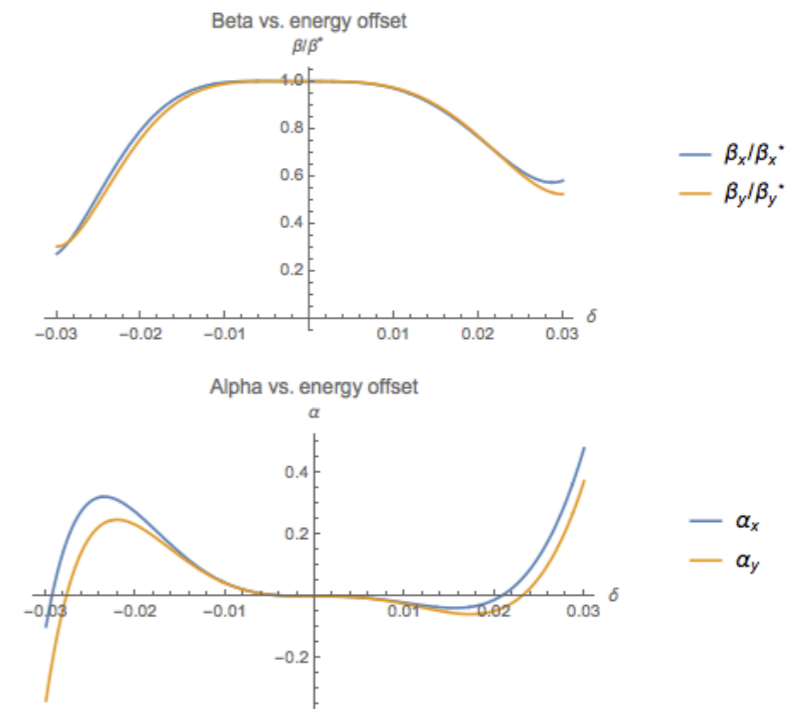


- 12 quads: cancel chromaticity to 2nd order.



For 0.5% energy spread: $(\Delta\epsilon_x/\epsilon_x, \Delta\epsilon_y/\epsilon_y) = (0.000186, 0.000186)$

For 1.% energy spread: $(\Delta\epsilon_x/\epsilon_x, \Delta\epsilon_y/\epsilon_y) = (0.0119, 0.0119)$

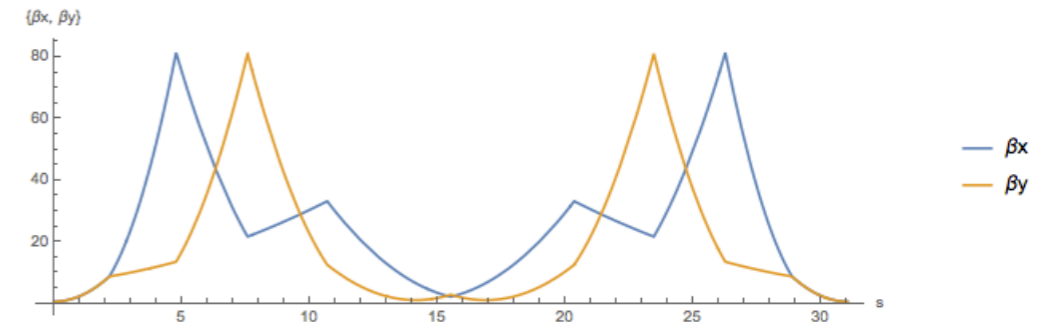


Length estimates and scalings

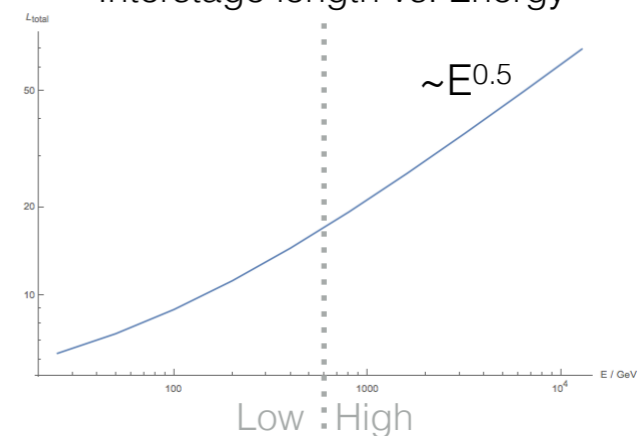
Any beta matching imposes: $L_{interstage} \sim \sqrt{E_{main}} \Rightarrow L_{collider} \sim E_{main}^{1.5}$

- High energy regime: **retains beta-shape**, constant emittance preservation, good synch. rad. scaling ($P_{SR} \sim E_{main}^2$)
- Low energy regime: complex lattices/many quads (high chrom. correction order), possibly use of sextupoles.
- Interstage **length estimate**: **~30 m @ 300 GeV**, (E-spread: 1% rms, dipole length: 1m, plasma ramp: 10x, quads: 150 T/m, emit. growth: ~1%, plasma density: 10^{16} cm^{-3})
- Note: work in progress.

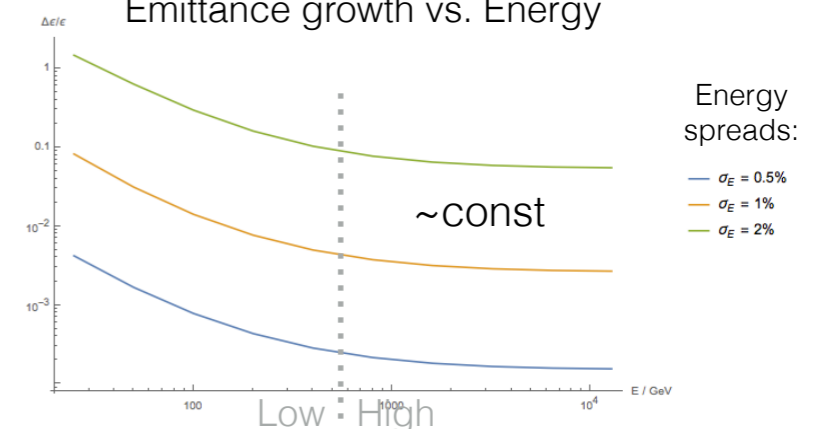
Current best solution (9 quads):



Interstage length vs. Energy



Emittance growth vs. Energy



Approximate scalings for high energy regime:

	Interstage length	Emittance growth
Energy	$E_{main}^{0.5}$	const
Quad strength	$g_{max}^{-0.5}$	$g_{max}^{-1.5}$
Energy spread (1st order chromaticity correction)	const	σ_E^4
Ramp magnification	const	$\prod p^{-3}$

Summary

- **Chromaticity is a big challenge** facing a PWFA interstage.
- Traditional chromaticity correction designs (using **sextupoles**) have **unfavourable energy scaling laws**.
- A new type of lattice:
Chromaticity-free quadrupole-only lattices have been developed (shorter, less SR, no non-linear terms).
- At high energies, **same optics solution applies to any energy** (length scales, emittance growth is constant)
- PWFA/LWFA **interstage and collider lengths will scale as:**
- Next up:
 - Integration of plasma density ramps
 - Emittance growth studies from SR/misalignment
 - Details of injection/extraction optics

$$L_{interstage} \sim \sqrt{E_{main}}$$

$$L_{collider} \sim E_{main}^{1.5}$$