

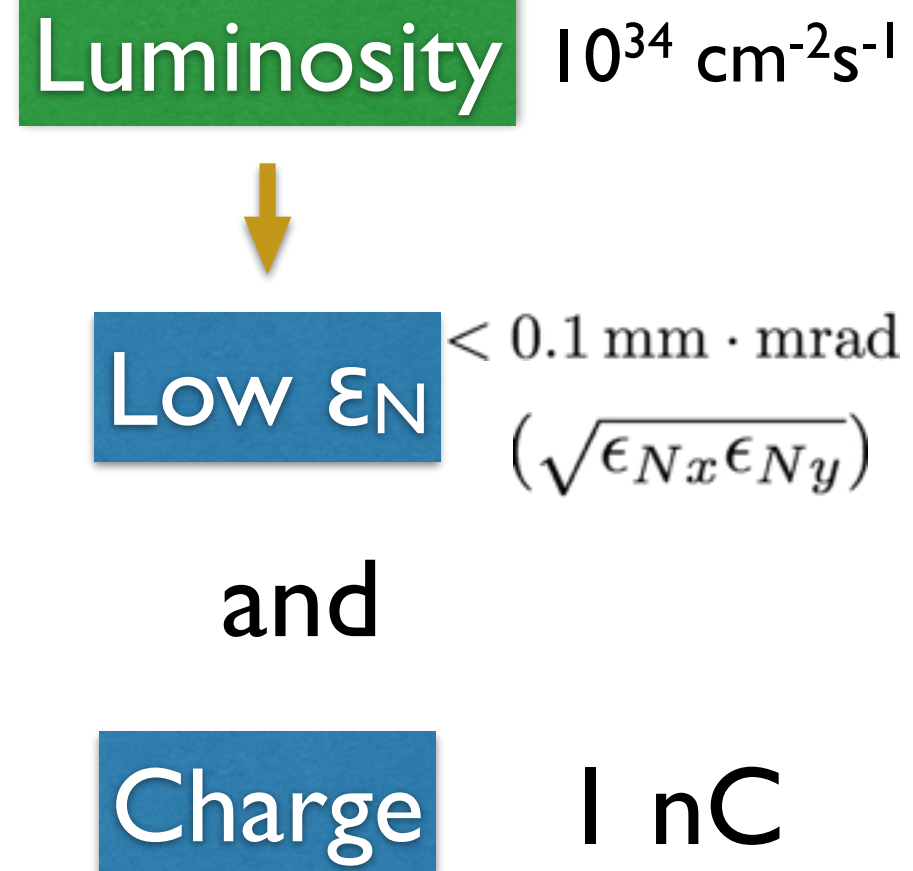
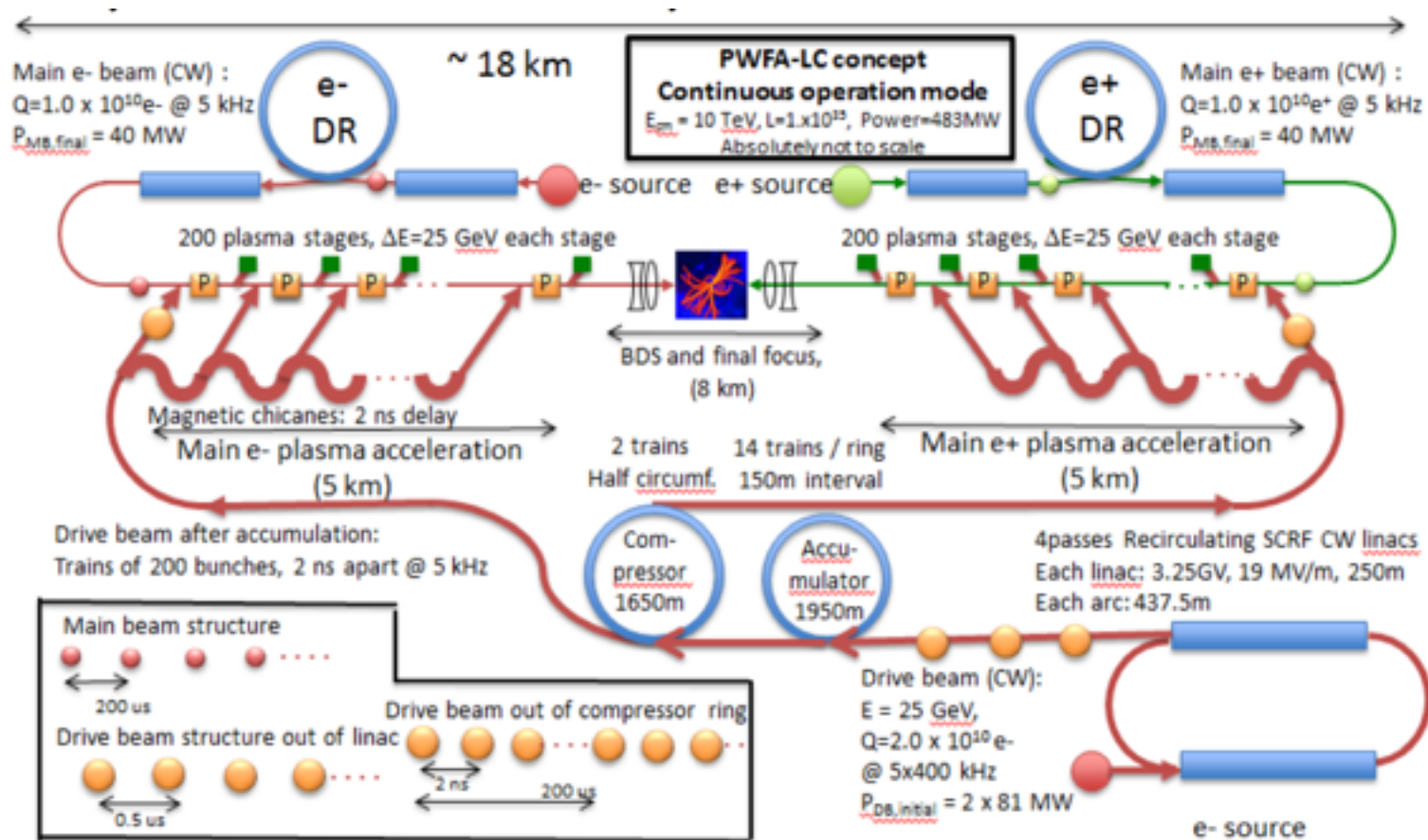
Plasma Ion Motion Induced Emittance Growth In Nonlinear Plasma Wake Field Accelerator

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PWFA Linear Collider

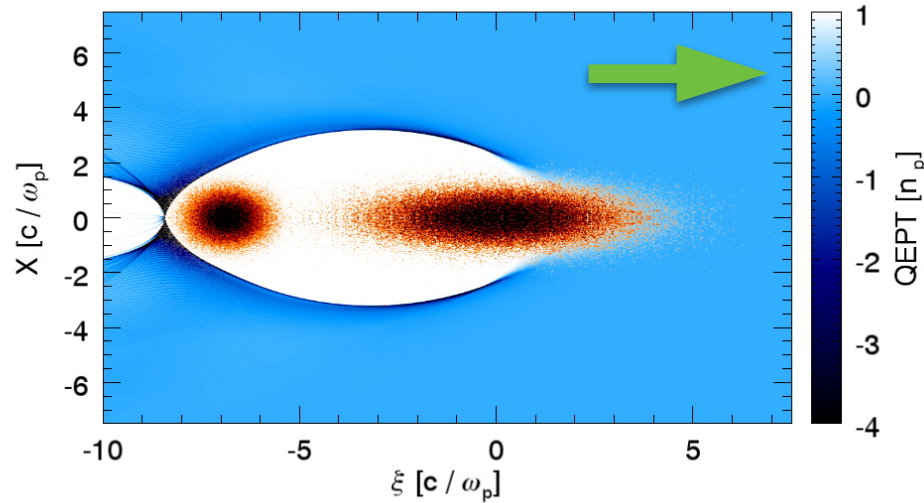
Conceptual design of a 10 TeV PWFA-LC



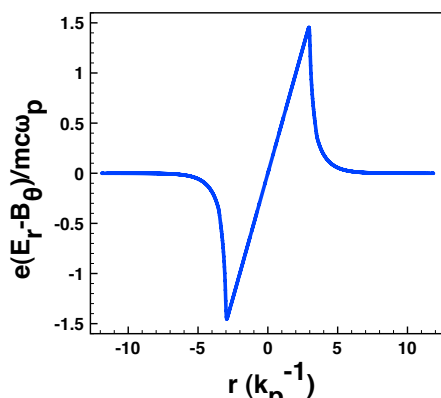
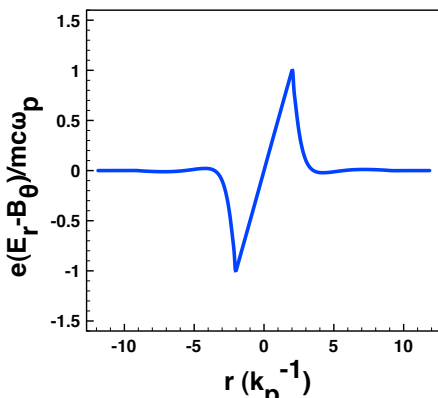
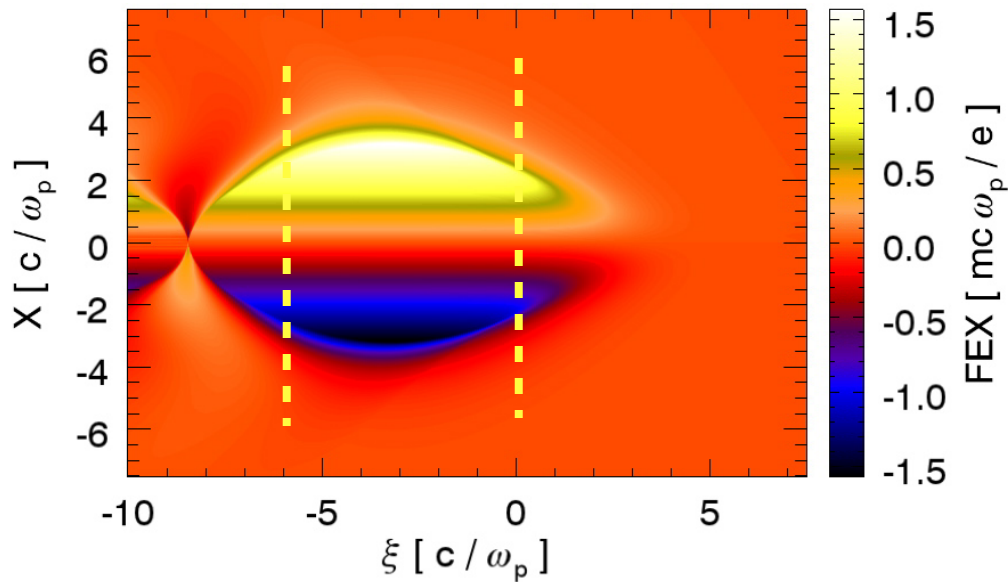
* E.Adli et al., IPAC 2014

Blowout PWFA

QEP



Focusing Field ($E_r - B_\theta$)



Matched spot size for an electron beam

$$\partial f / \partial t = 0$$

$$\sigma_{r,\text{matched}} [\mu\text{m}] = 1.304 \left(\frac{\epsilon_N^2 [\text{mm} \cdot \text{mrad}]}{E [\text{GeV}] n_p [10^{16} \text{cm}^{-3}]} \right)^{\frac{1}{4}}$$

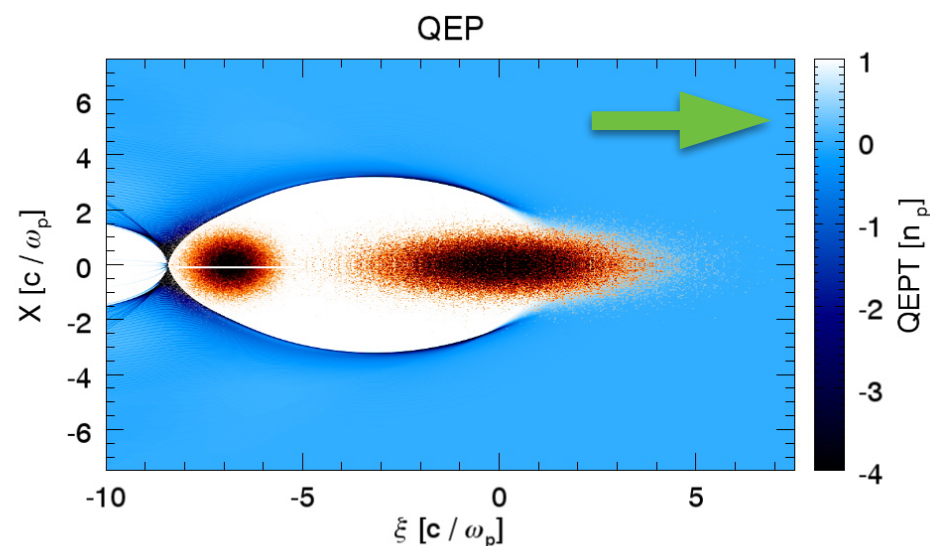
$$n_p = 1.0 \times 10^{17} \text{ cm}^{-3}$$

$$E = 25 \text{ GeV},$$

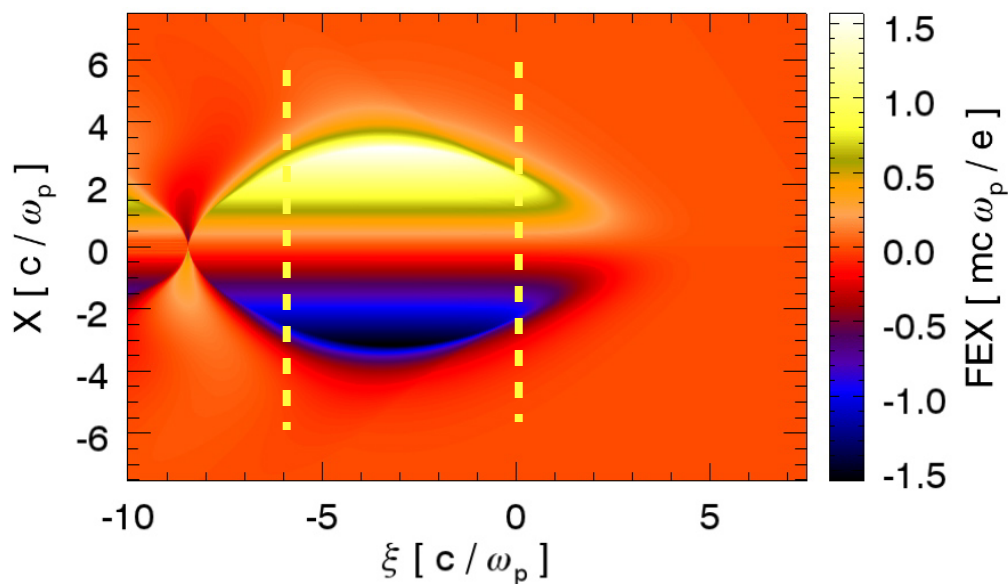
$$\epsilon_N = 0.05 \text{ mm} \cdot \text{mrad}$$

$$\sigma_r = 73.3 \text{ nm}$$

Blowout PWFA



Focusing Field ($E_r - B_\theta$)



$\sigma_r = 73.3 \text{ nm}$ \rightarrow Ions collapse!

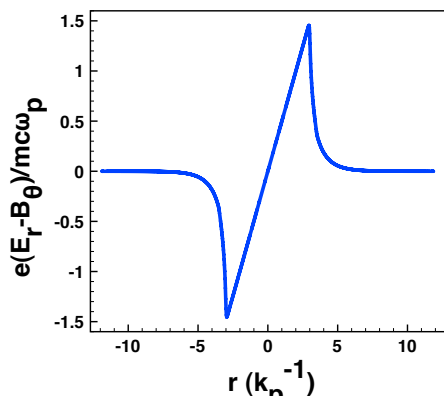
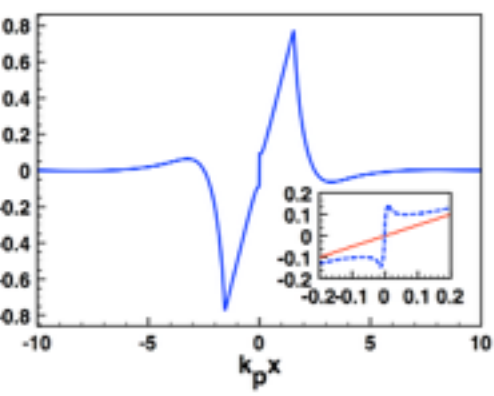


$$\Omega_b = \sqrt{4\pi n_b e^2 / m_i} , \Omega_b \sigma_z / c \gg 1$$

* W. An et al. Phys. Rev. Lett. 118, 244801, 2017

$$\sigma_z \approx c / \omega_p , \frac{n_b}{n_p} \gg \frac{m_{ion}}{m_e}$$

* J. B. Rosenzweig et al. Phys. Rev. Lett. 95, 195002, 2005



$\sigma_r = 0.1 \mu\text{m}, \sigma_z = 10.0 \mu\text{m}, N = 1.0 \times 10^{10}$
 $n_p = 1.0 \times 10^{17} \text{ cm}^{-3}$
 $n_b/n_p = 63500 \gg m_{ion}/m_e = 1836(\text{H}), 12700(\text{Li})$

* J. B. Rosenzweig et al. Phys. Rev. Lett. 95, 195002, 2005

TABLE I. Beam and plasma parameters for linear collider afterburner, derived from Ref. [7].

N_b (drive, accelerating)	$1.5 \times 10^{10}, .5 \times 10^{10}$
rms bunch length σ_z	$35 \mu\text{m}$
γ (drive, accelerating)	$\leq 1 \times 10^6, \leq 2 \times 10^6$
Accelerated beam $\varepsilon_{n,(x,y)}$	$4 \times 10^{-6}, 9.6 \times 10^{-6} \text{ m rad}$
Drive beam $\varepsilon_{n,x}$	$6.2 \times 10^{-7} \text{ m rad}$
Initial ion (electron) density n_0	$0.9 \times 10^{16} \text{ cm}^{-3}$
Ion charge state Z	1 (hydrogen)
Matched β function β_{eq}	3.1 cm
Normalized beam density n_b/n_0	1.5×10^5

$$n_p = 1.0 \times 10^{16} \text{ cm}^{-3} \text{ Hydrogen}$$

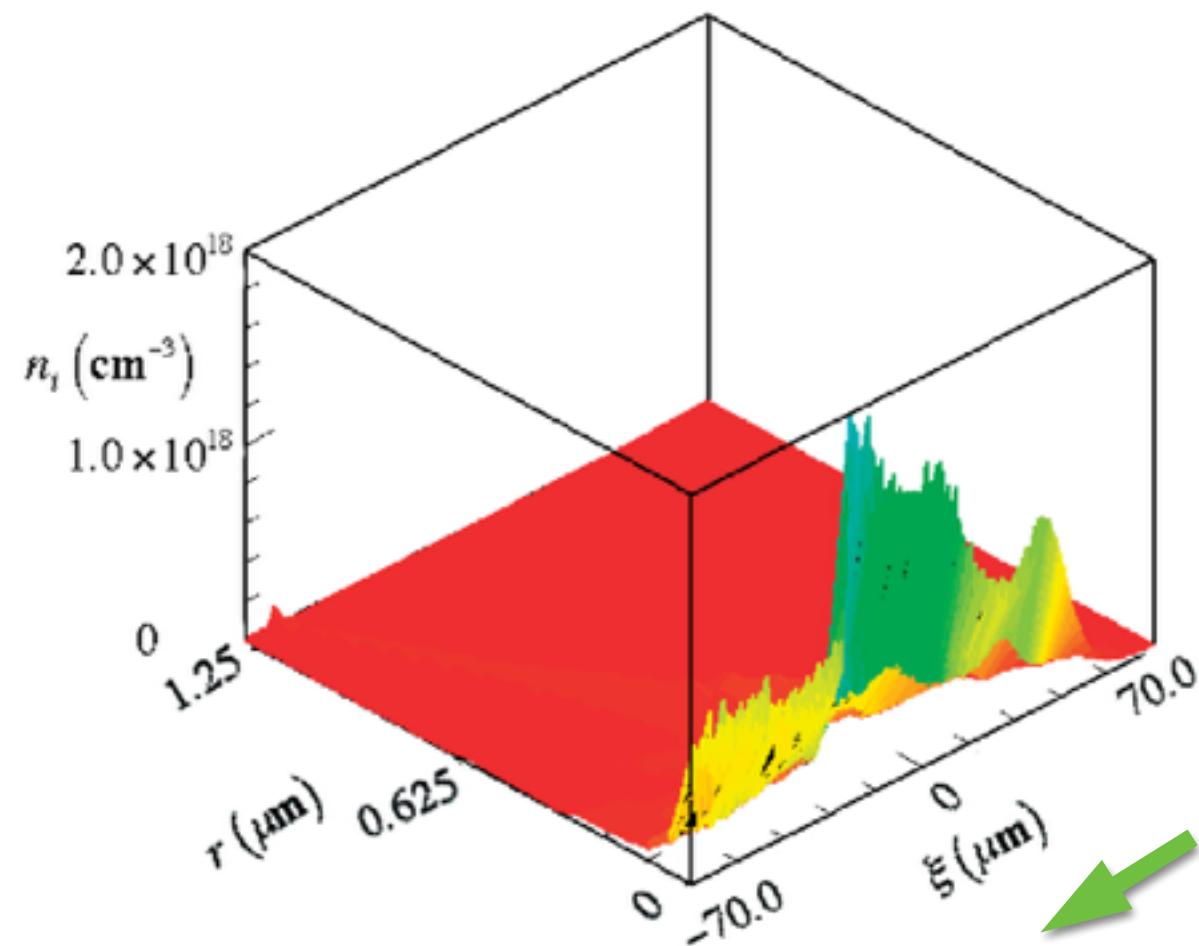


FIG. 1 (color). Surface plot of ion density distribution in (ζ, r) , as simulated by OOPIC for drive beam conditions of Table I.

The Ion profile: $1 + A \exp(-r^2/2\sigma^2)$

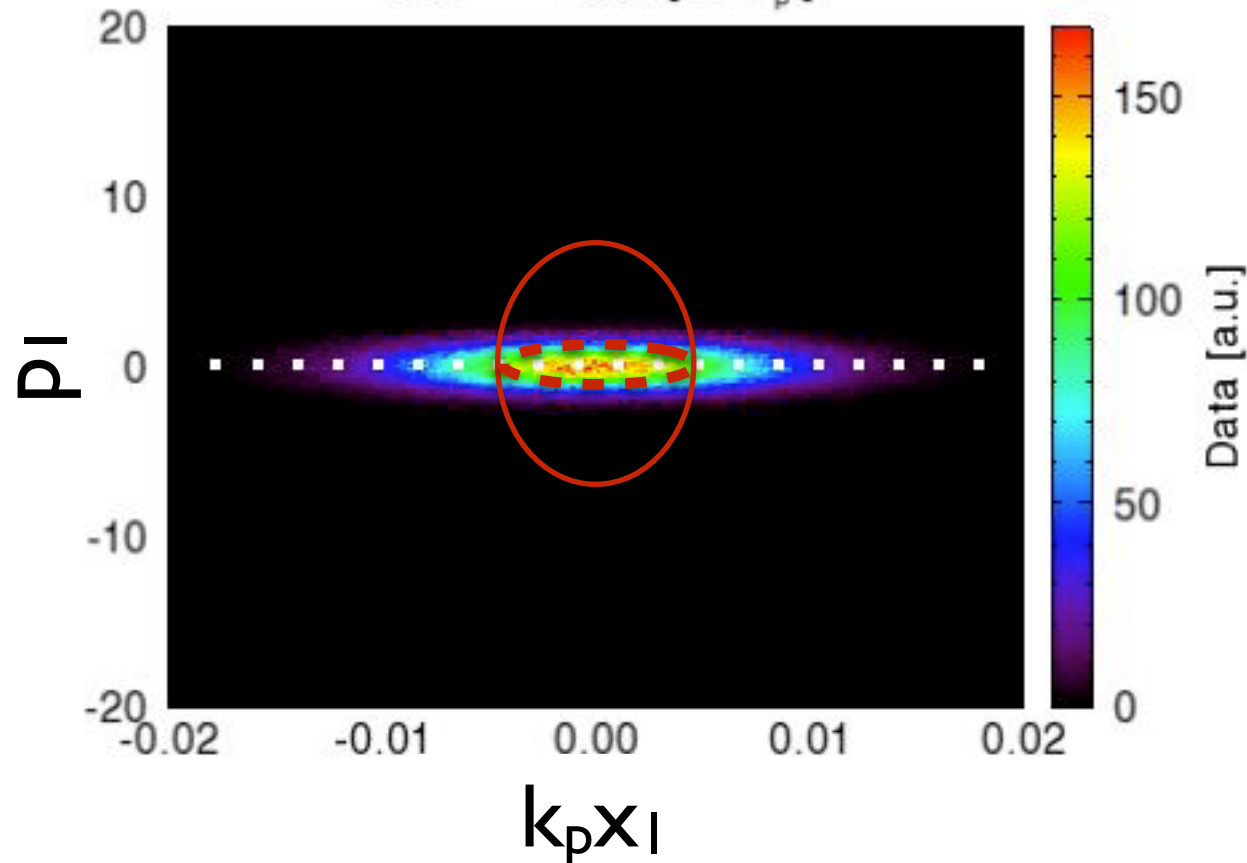
The focusing force: $k_p r/2 + A(k_p \sigma)^2 \frac{1 - \exp(-r^2/2\sigma^2)}{k_p r}$

Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$, $\sigma_z = 10.0 \mu\text{m}$, $N = 1.0 \times 10^{10}$

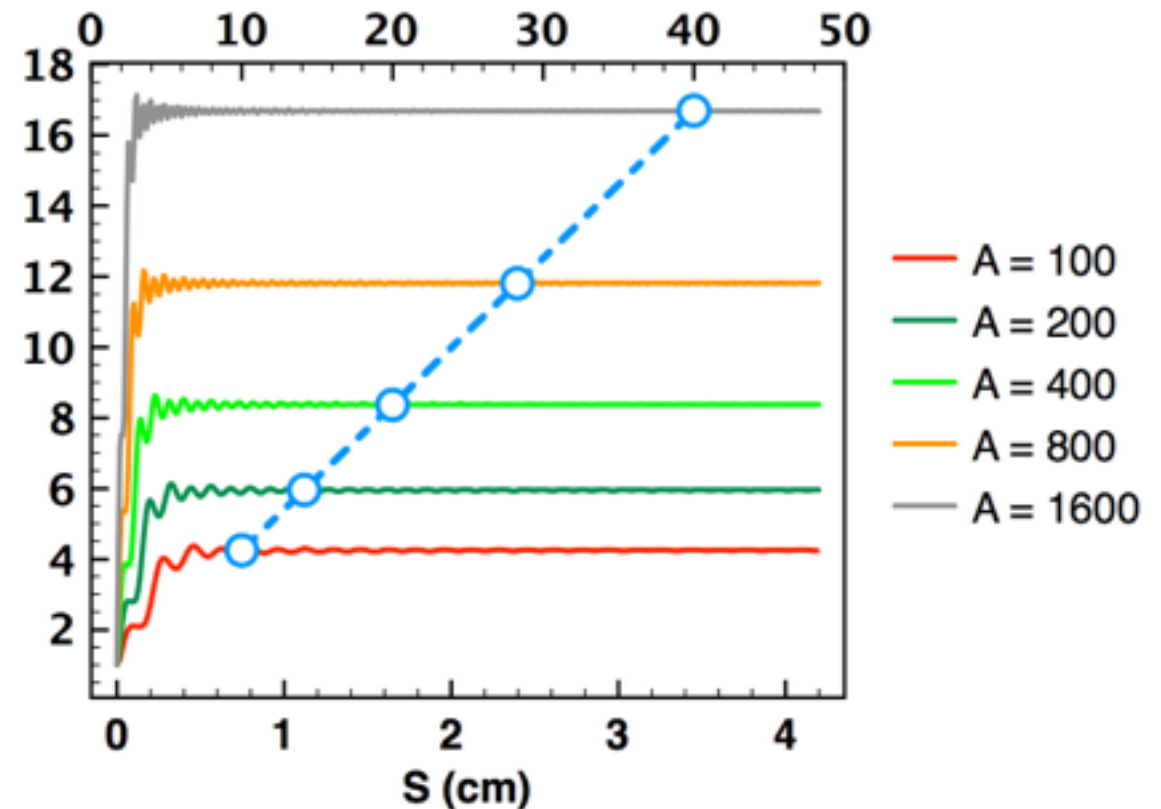
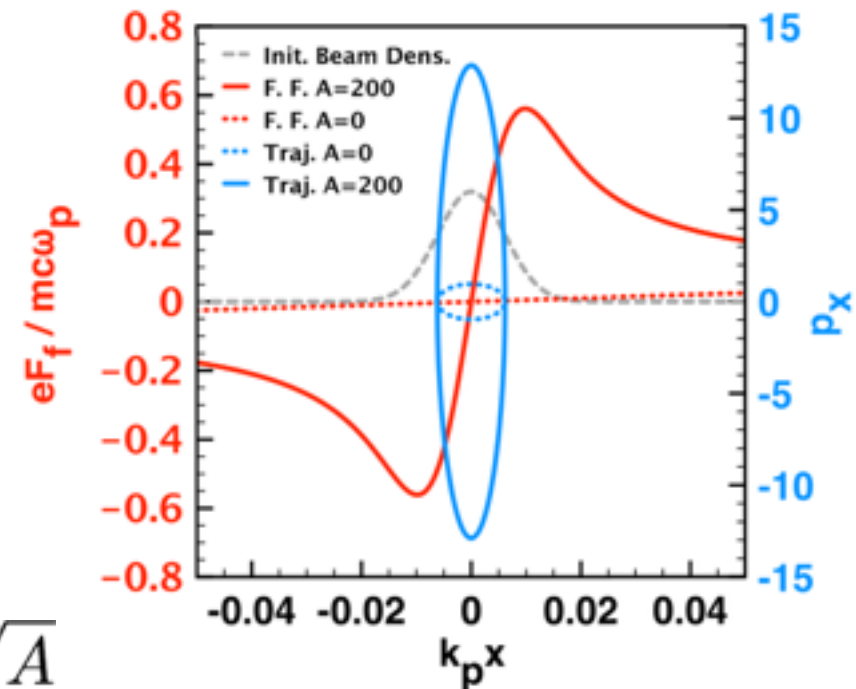
$\epsilon_N = 0.093 \text{ mm}\cdot\text{mrad}$

Phase Space of $p|x|$

Time = 0.00 [$1/\omega_p$]



Ion Profile: $A=100.0$, $\sigma = 0.1 \mu\text{m}$



$$\epsilon_{Nf}/\epsilon_{N0} \propto p_{\text{max}} \propto A^{0.5}$$

The Ion profile: $1 + A \exp(-r^2/2\sigma^2)$

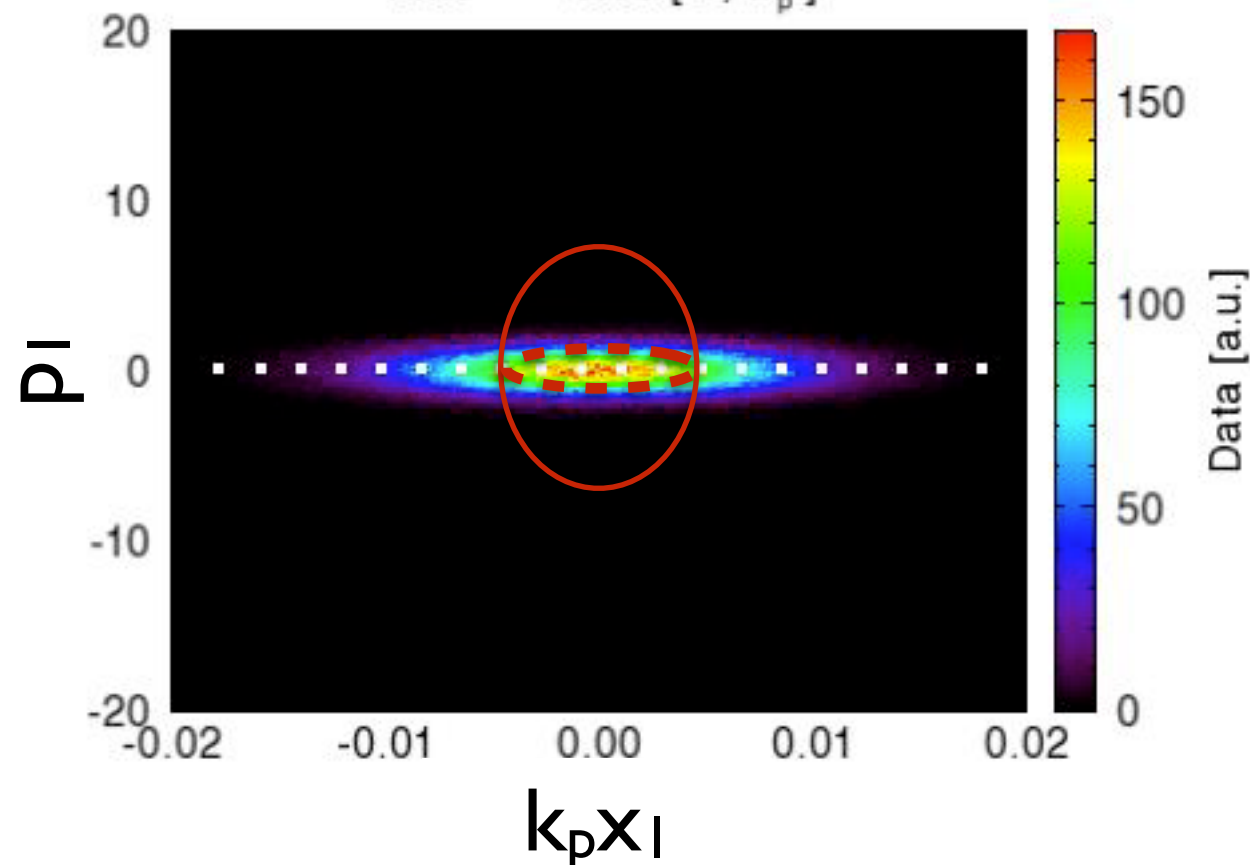
The focusing force: $k_p r/2 + A(k_p \sigma)^2 \frac{1 - \exp(-r^2/2\sigma^2)}{k_p r}$

Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$, $\sigma_z = 10.0 \mu\text{m}$, $N = 1.0 \times 10^{10}$

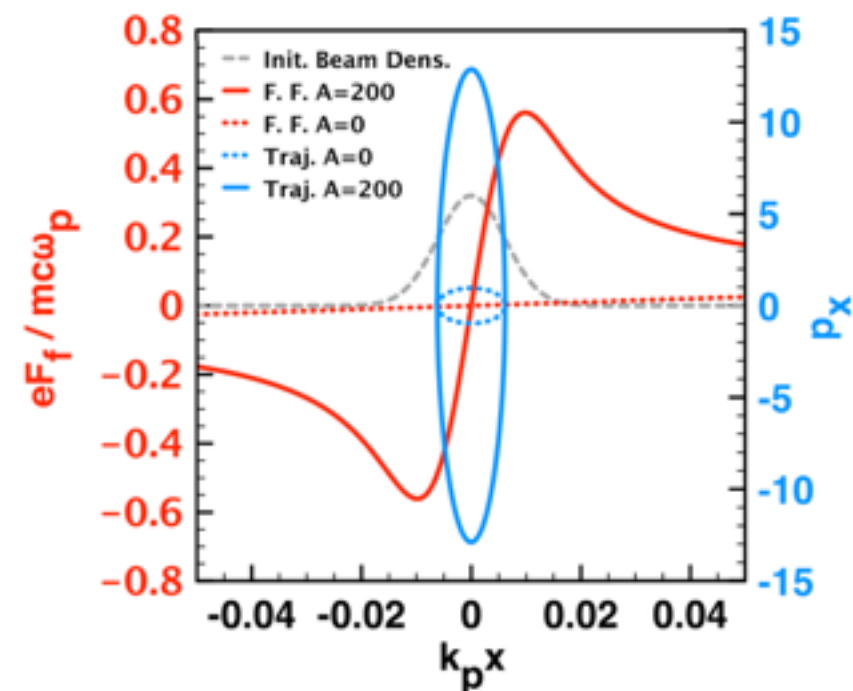
$\epsilon_N = 0.093 \text{ mm}\cdot\text{mrad}$

Phase Space of $p|x|$

Time = 0.00 [1/ω_p]



Ion Profile: $A=100.0$, $\sigma = 0.1 \mu\text{m}$



$$\epsilon_N = \gamma \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2}$$

$$\epsilon'_N = \frac{\gamma}{2\epsilon_N} (\langle x^2 \rangle \langle x'x'' \rangle - \langle xx' \rangle \langle xx'' \rangle)$$

$$\text{If } x'' = cx, \epsilon'_N = 0$$

* J. B. Rosenzweig et al. Phys. Rev. Lett. 95, 195002, 2005

* R. Gholizadeh et al. Phys. Rev. Lett. 104, 155001, 2010

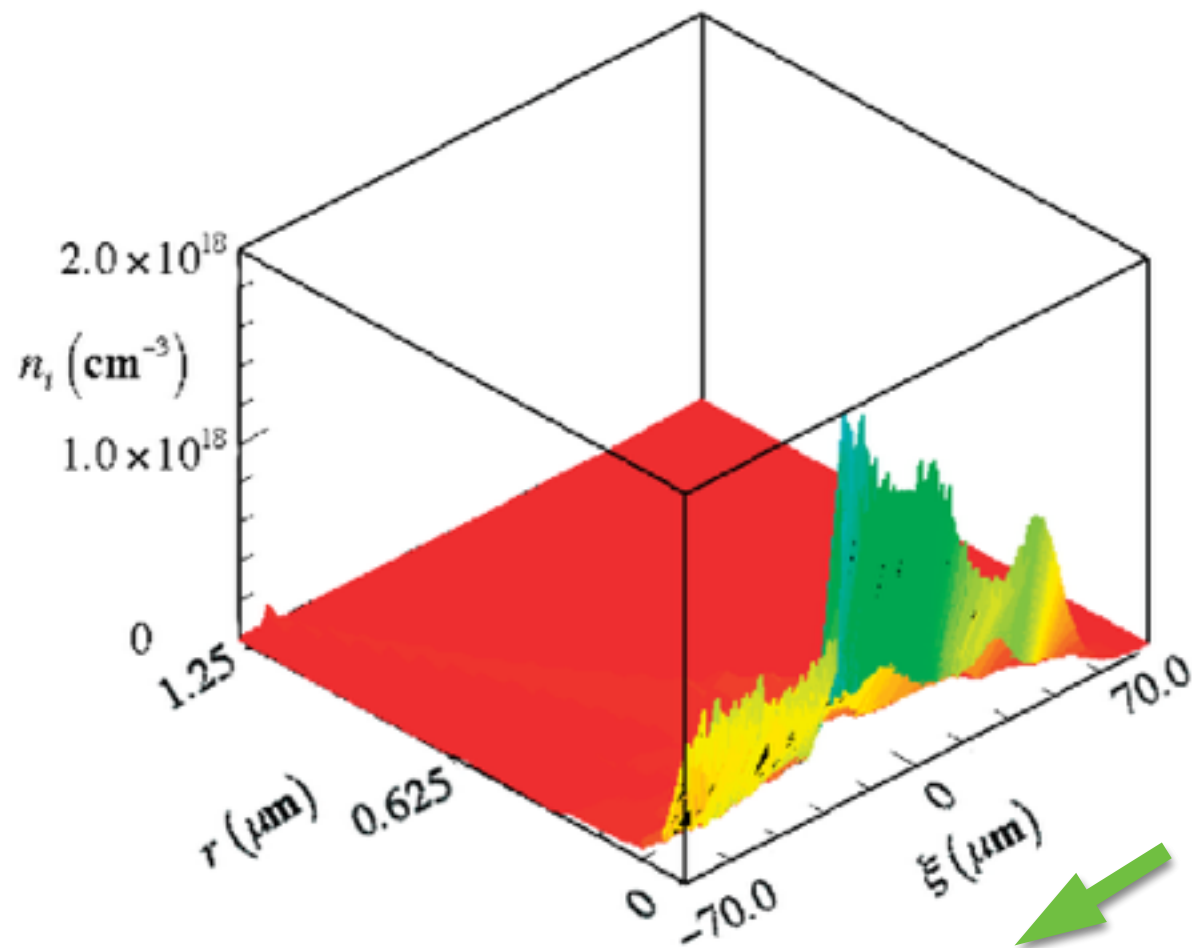
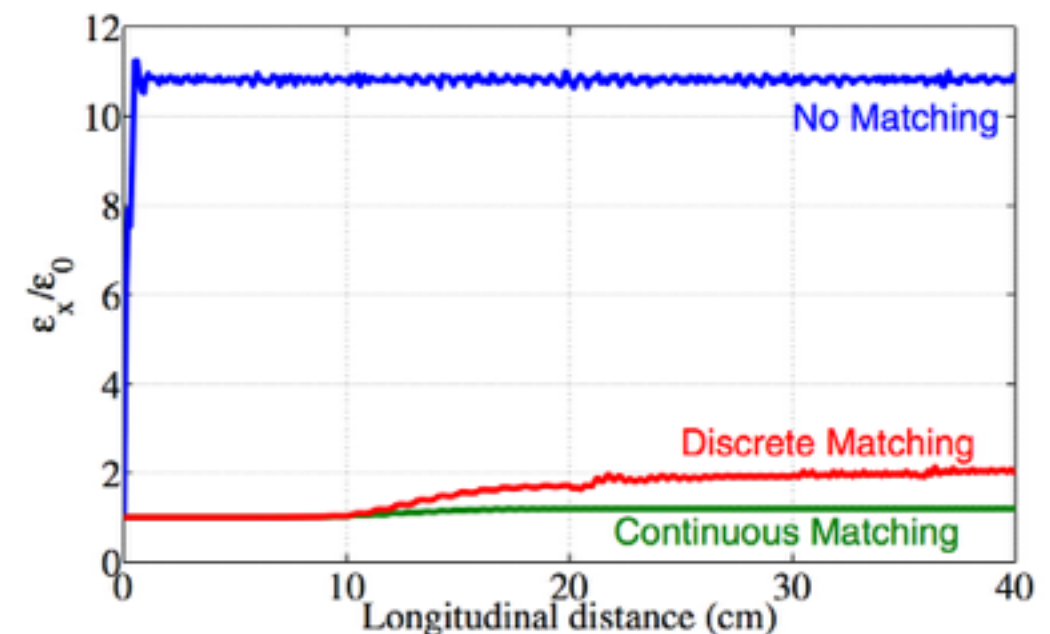
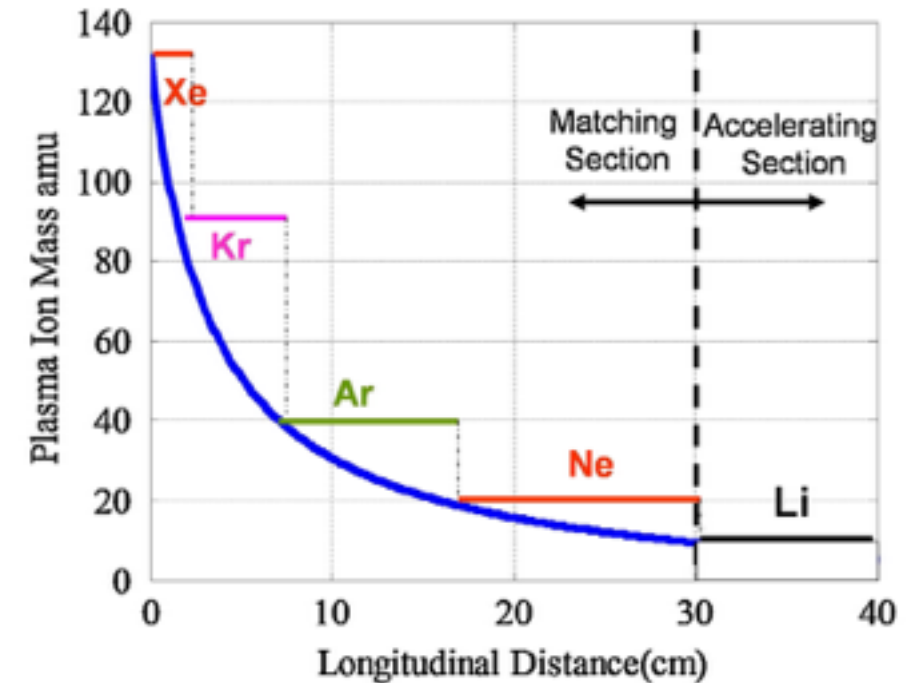


FIG. 1 (color). Surface plot of ion density distribution in (ζ, r) , as simulated by OOPIC for drive beam conditions of Table I.

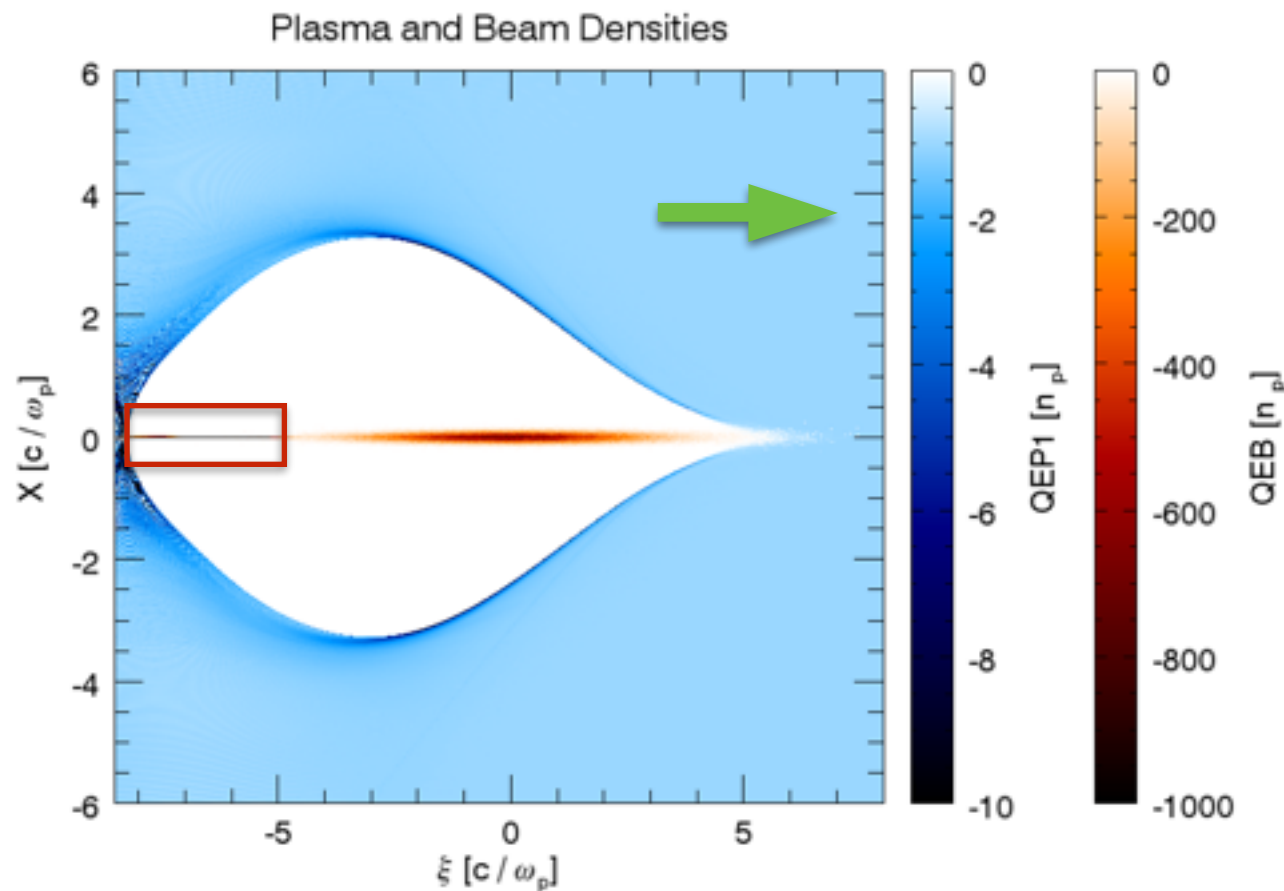
$$d\varepsilon_{n,x}/dz \approx 6 \times 10^{-4} \text{ m rad/m}$$



3D Simulation Big Challenge

$$\sigma_r = 0.1 \mu\text{m}, \sigma_z = 10.0 \mu\text{m}, N = 1.0 \times 10^{10}$$

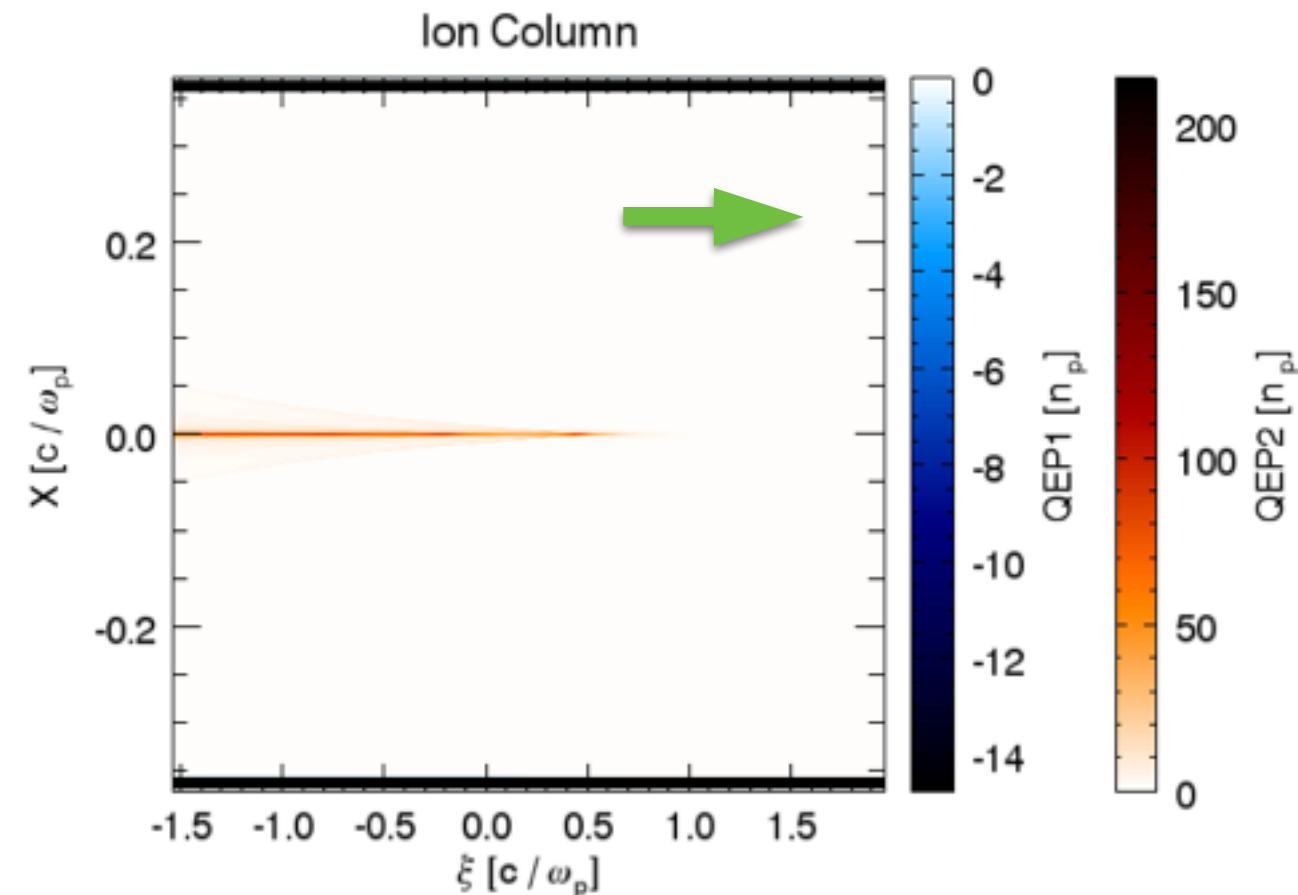
$$n_b/n_p = 63500 \gg m_{\text{ion}}/m_e = 1836$$



400 μm x 400 μm x 300 μm Box

8192 x 8192 x 1024 Cells

$$\Delta_{\perp} \approx 50 \text{ nm}$$



12 μm x 12 μm x 60 μm Box

4096 x 4096 x 512 Cells

$$\Delta_{\perp} \approx 3 \text{ nm}$$

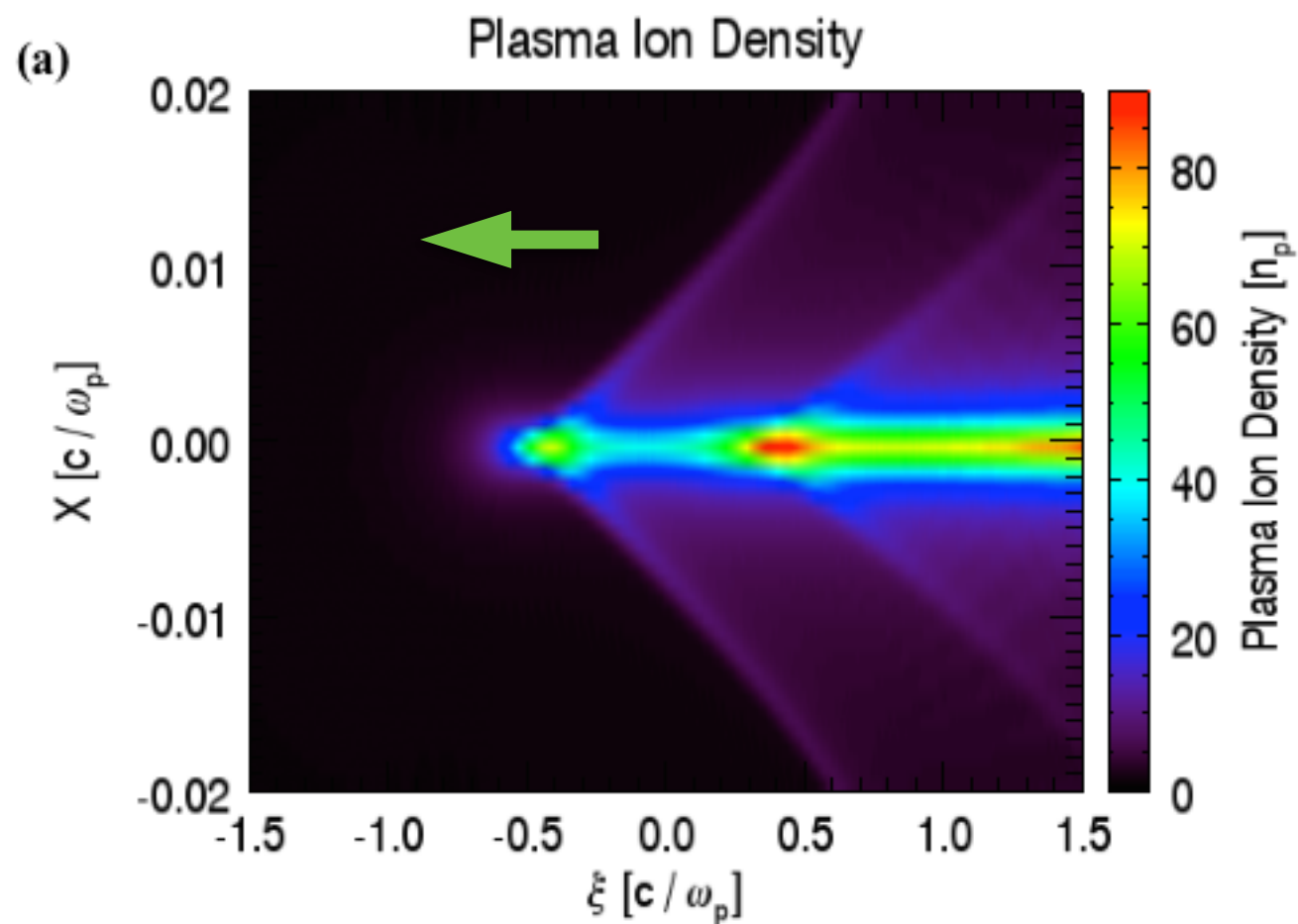
400 μm x 400 μm x 313 μm Box

12 μm x 12 μm x 60 μm Box

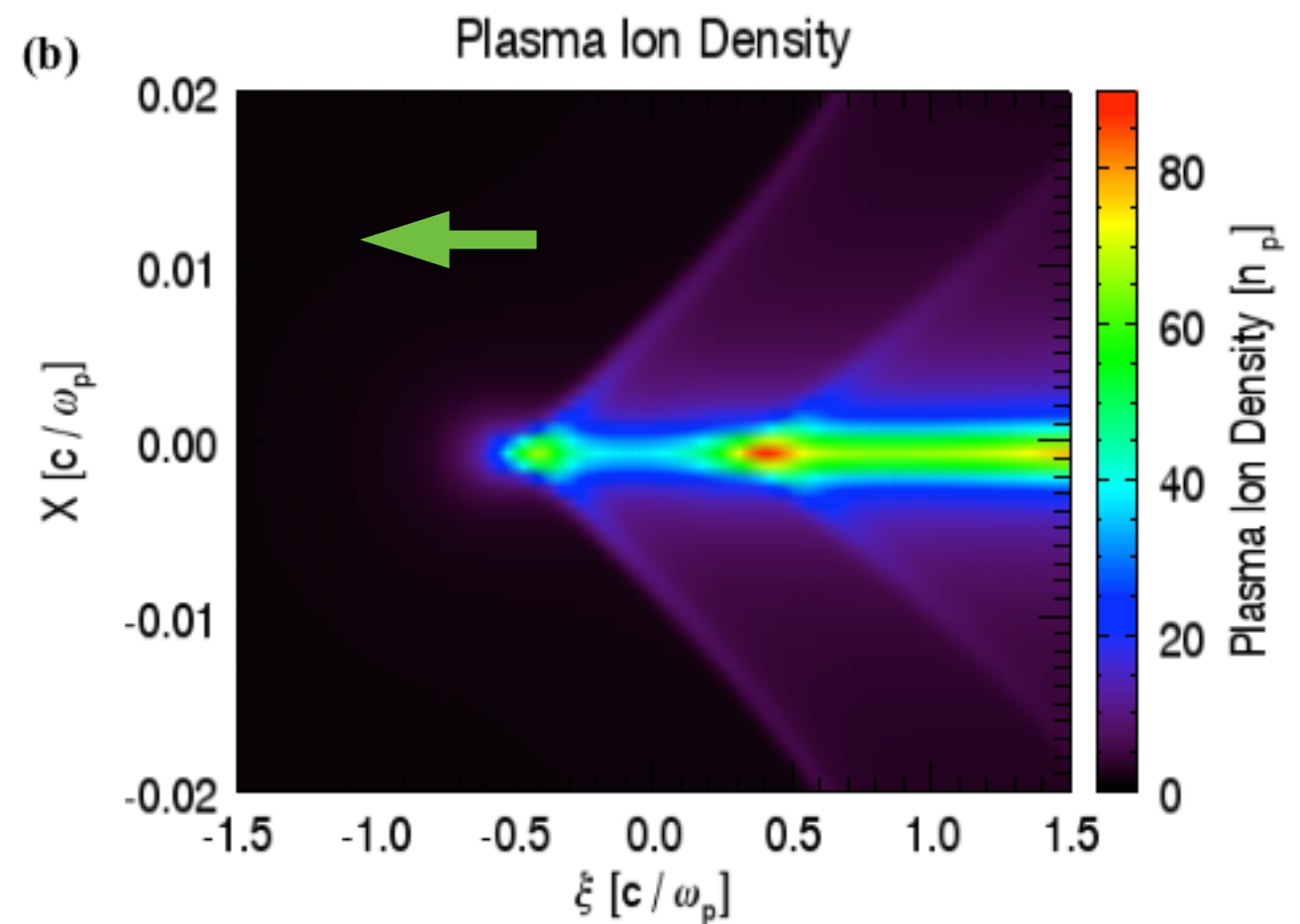
$\Delta_{\perp} \approx 12 \text{ nm}$

32768 x 32768 x 2693 Cells

1024 x 1024 x 512 Cells

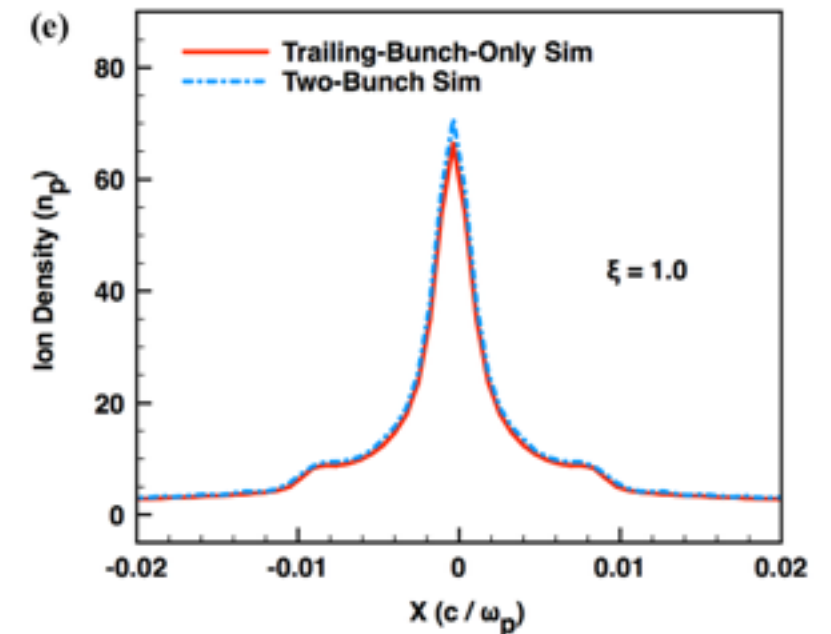
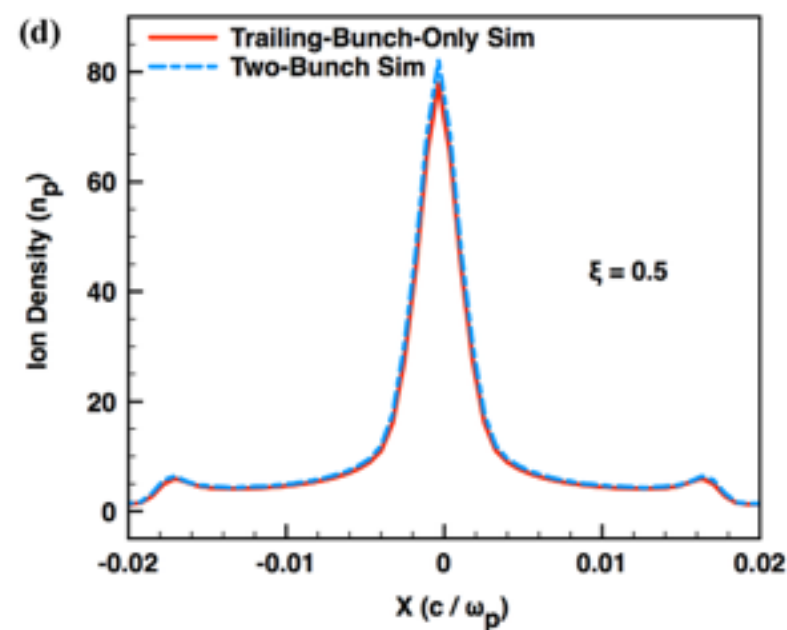
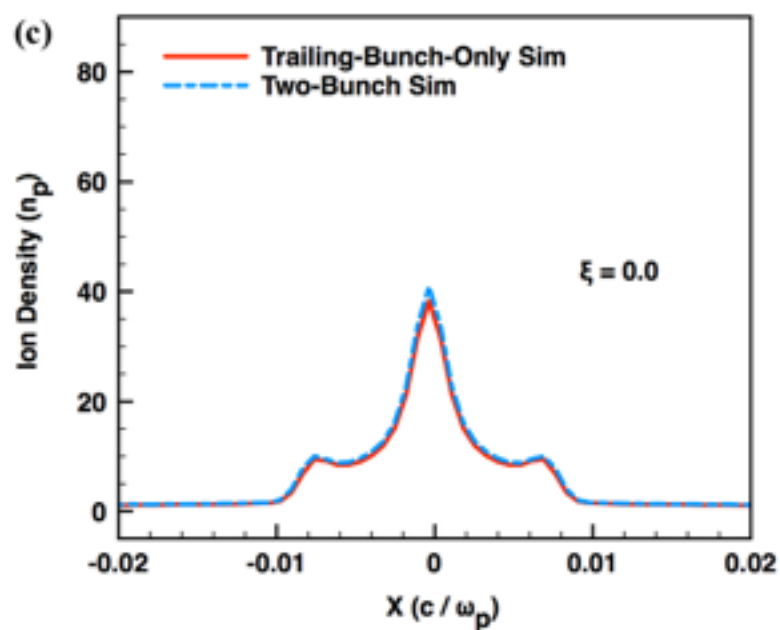
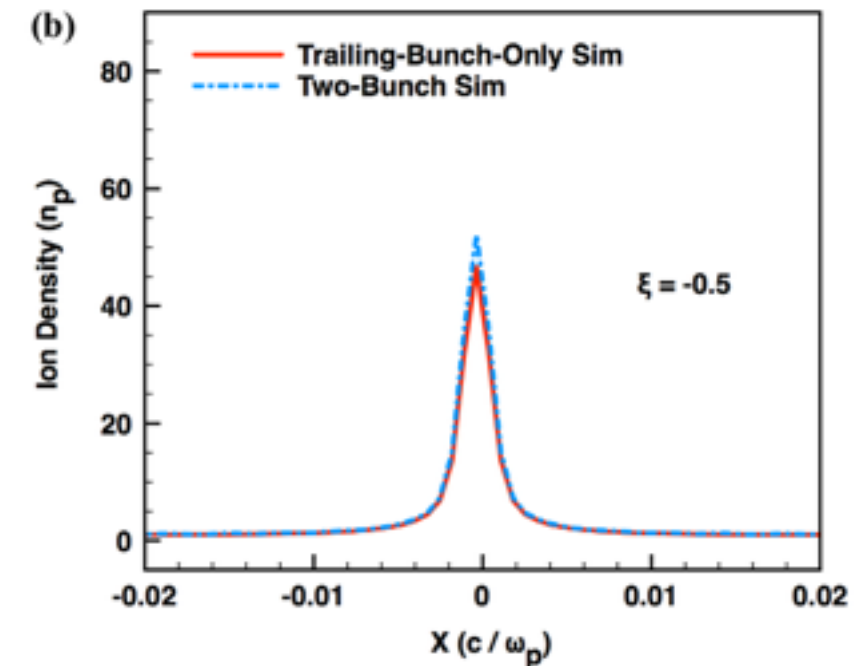
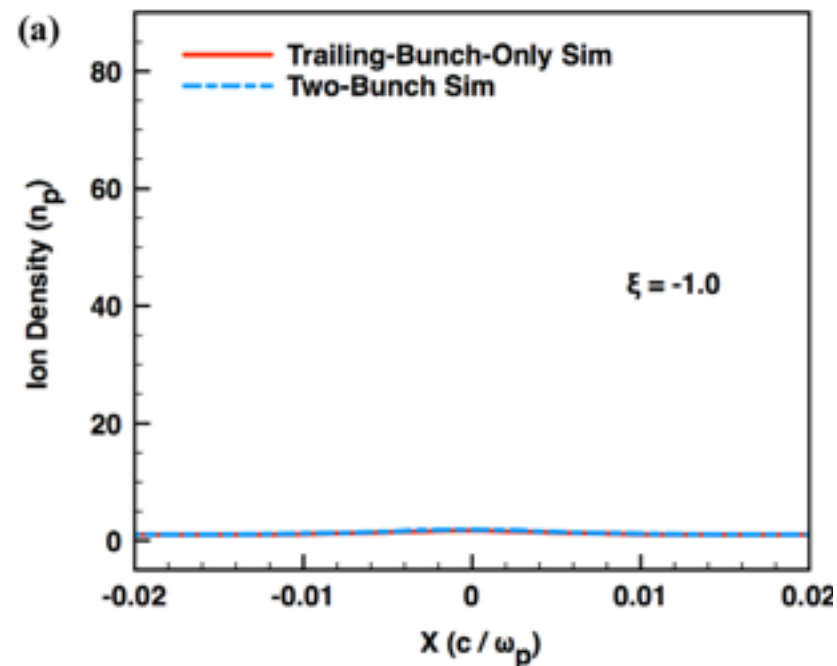
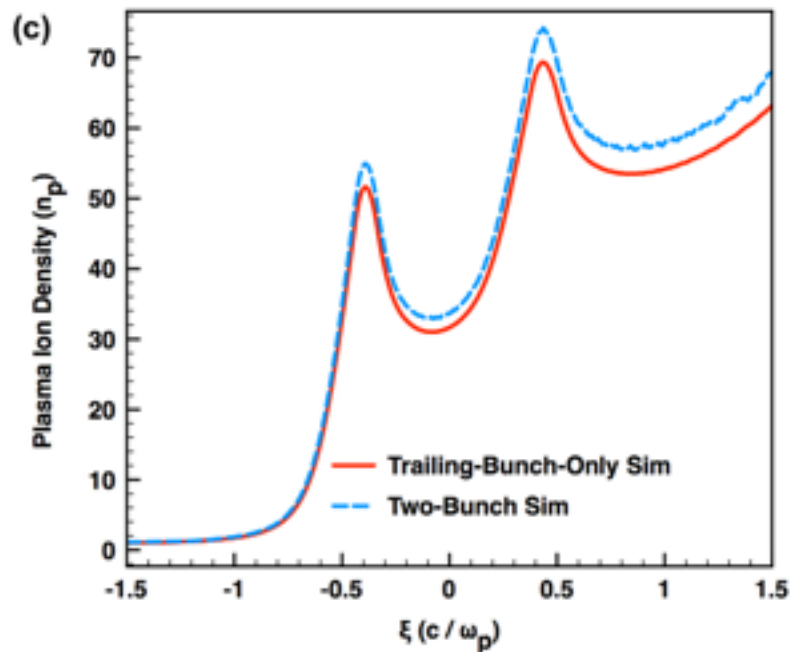


Two-Bunch



Trailing-Bunch-Only

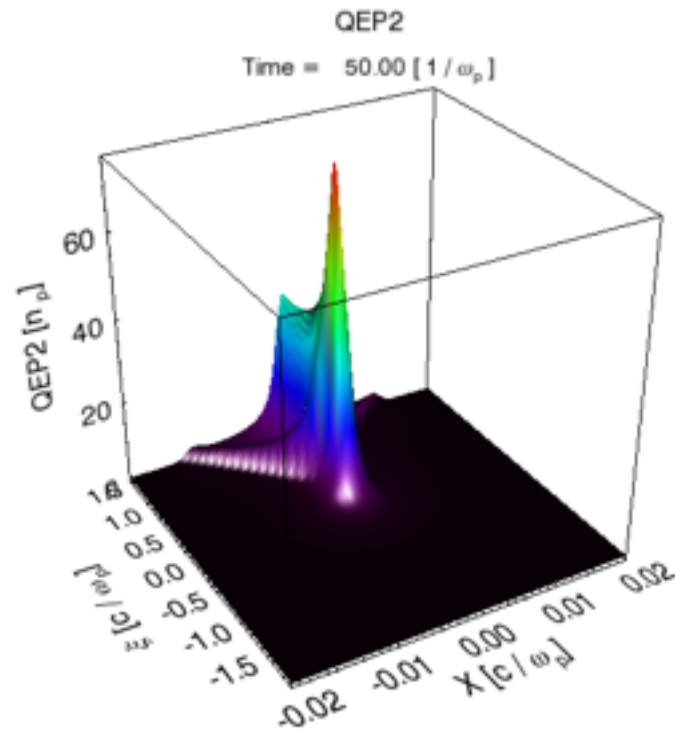
Comparison of Longitudinal and Transverse Lineouts



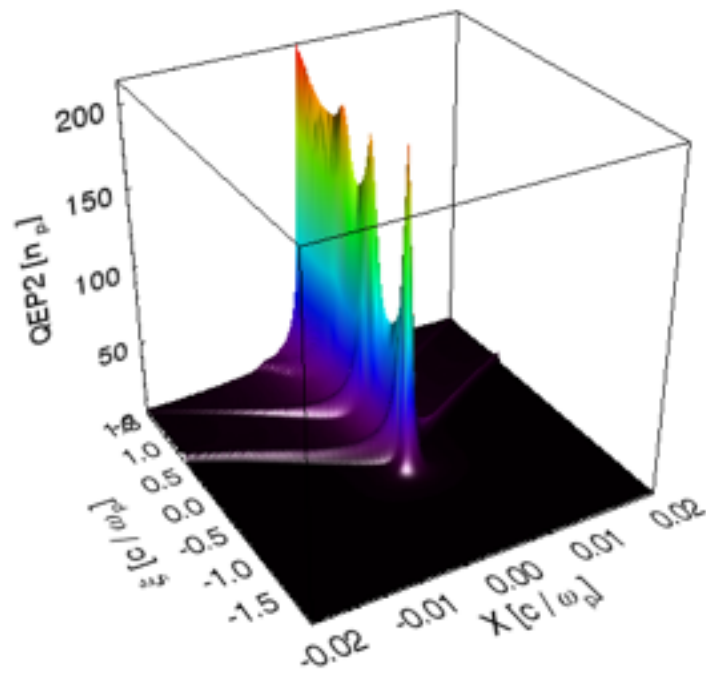
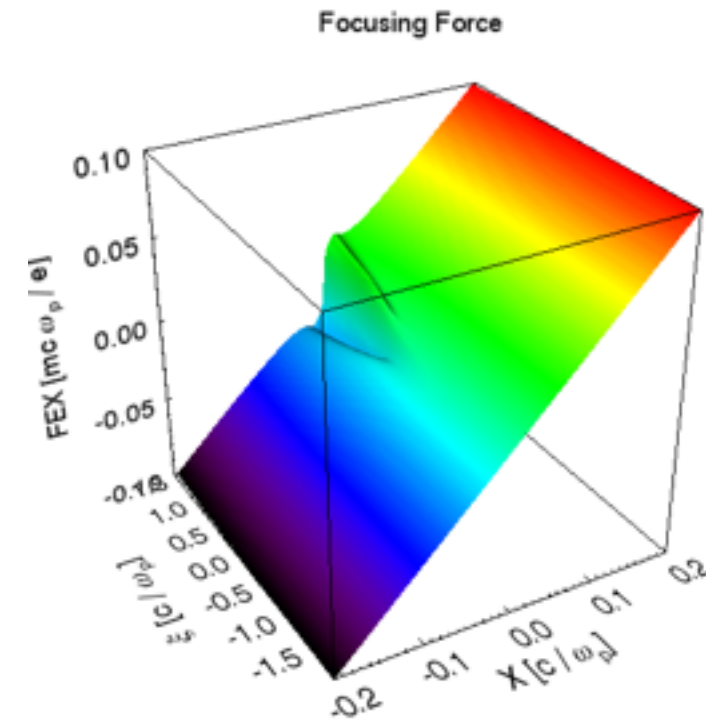
Drive Beam : $\sigma_r = 1 \mu\text{m}$, $\sigma_z = 30.0 \mu\text{m}$, $N_1 = 3.0 \times 10^{10}$, $\epsilon = 10 \text{ mm}\cdot\text{mrad}$

Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$ ($0.006 k_p^{-1}$) , $\sigma_z = 10.0 \mu\text{m}$, $N_2 = 1.0 \times 10^{10}$, $\epsilon = 0.093 \text{ mm}\cdot\text{mrad}$

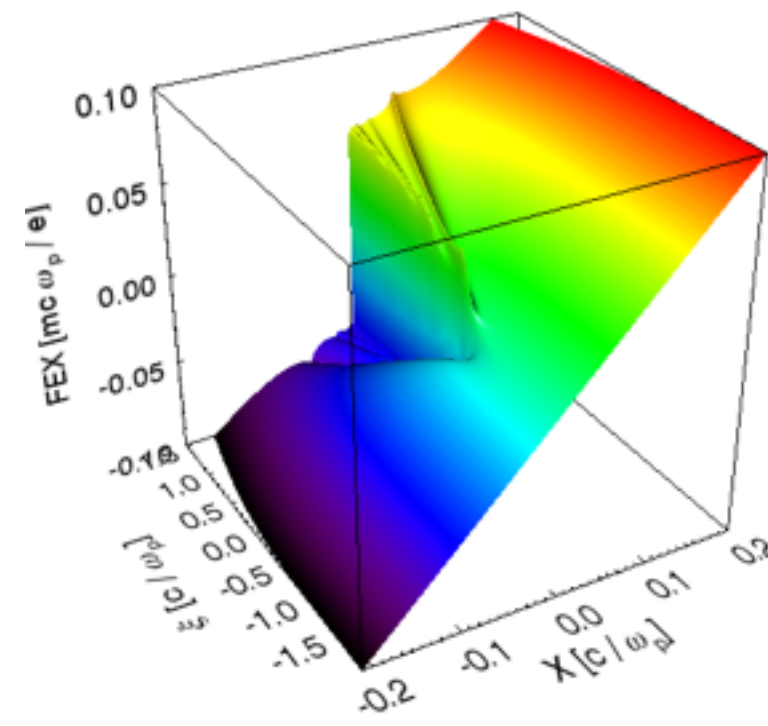
Distance between two beams : $115 \mu\text{m}$; Plasma Density : $1.0 \times 10^{17} \text{ cm}^{-3}$



Li^+

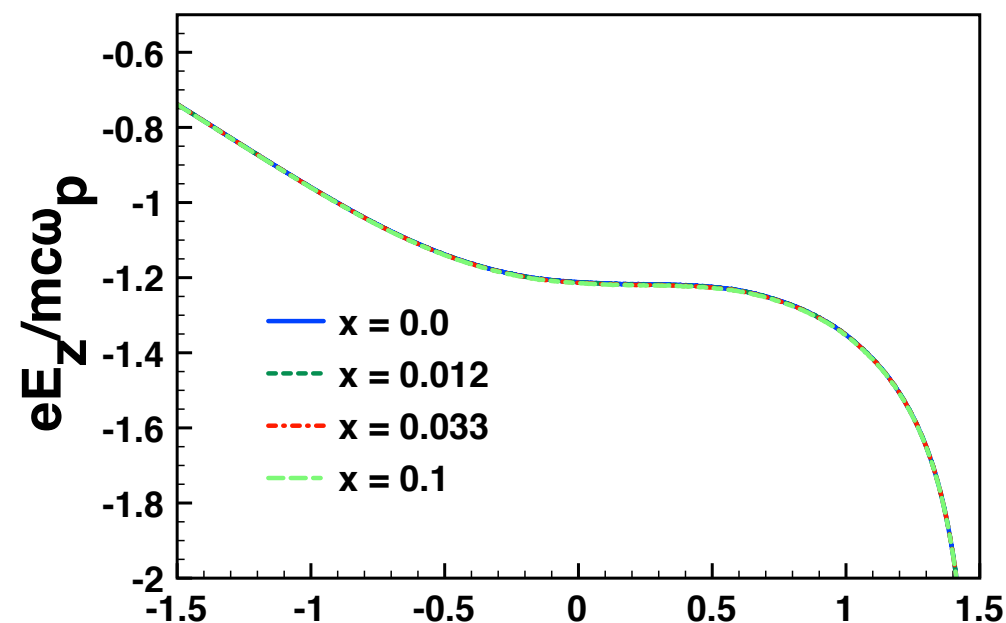
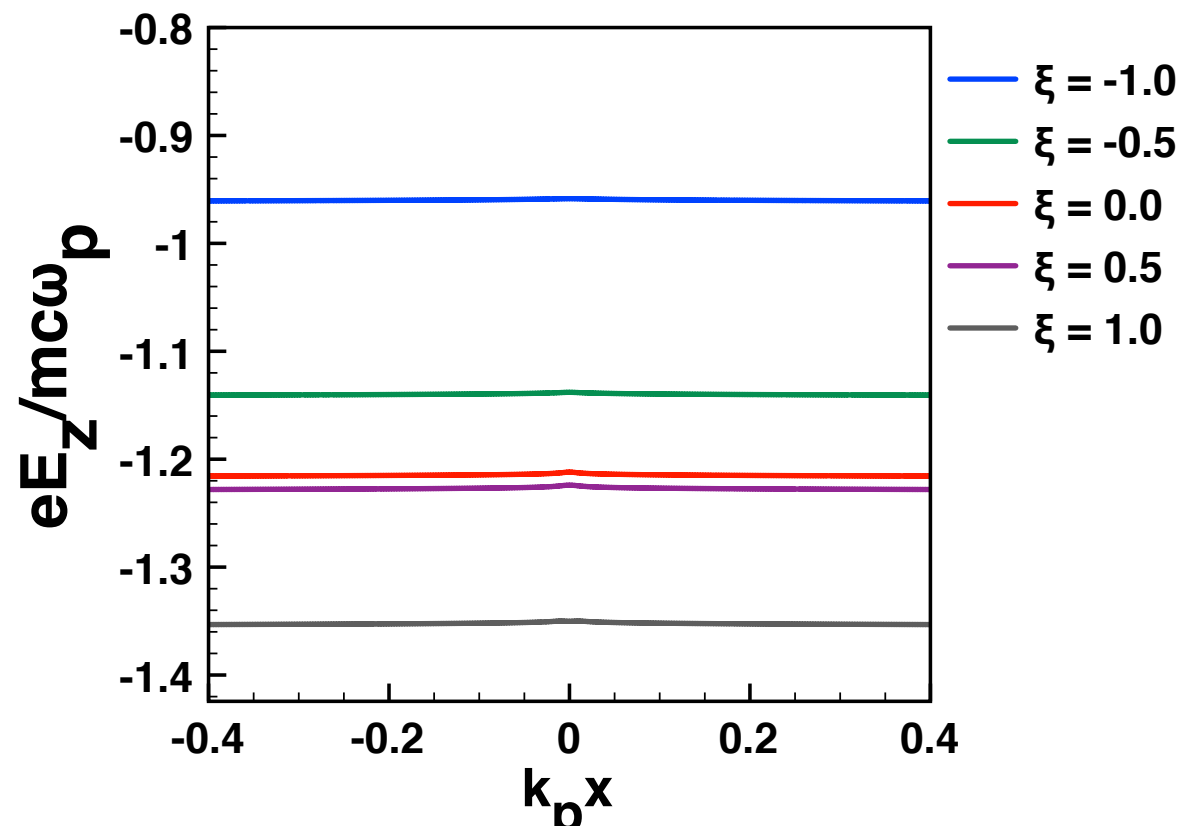
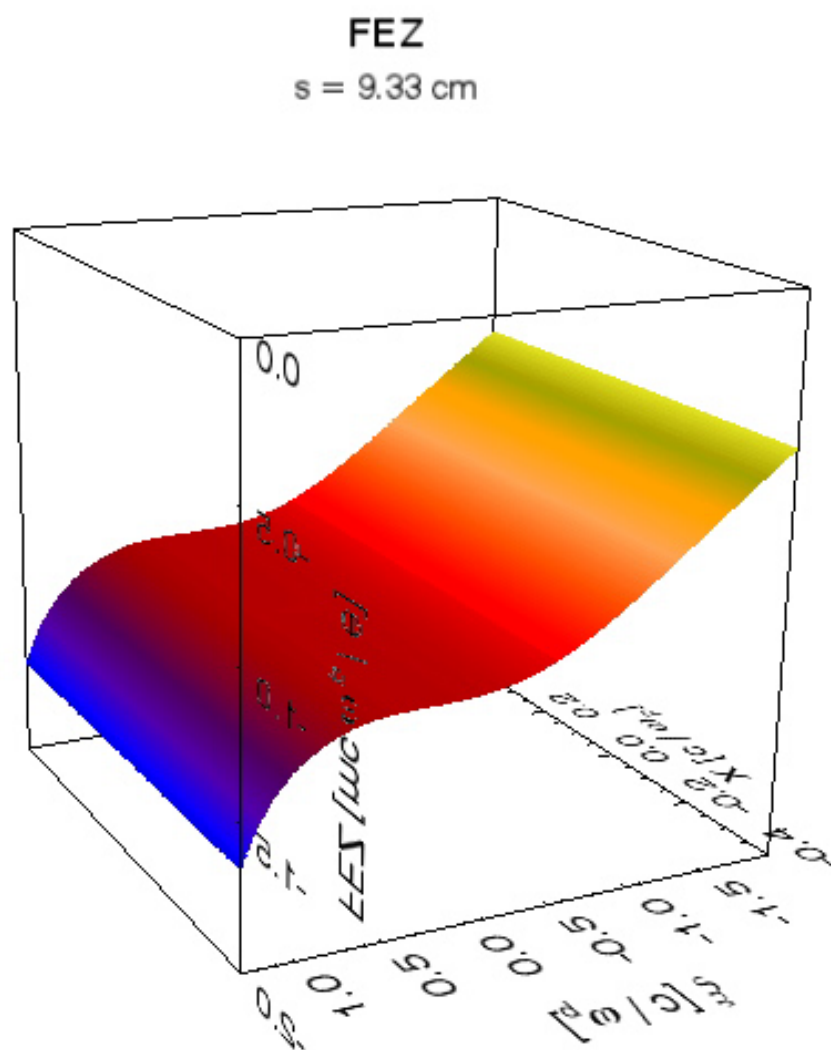


H^+



The accelerating field of the plasma wake

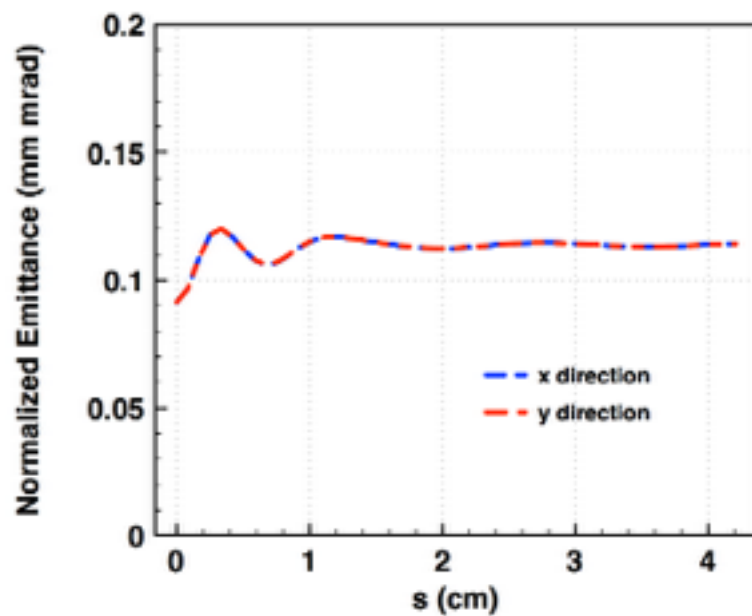
Drive Beam : $\sigma_r = 1 \mu\text{m}$, $\sigma_z = 30.0 \mu\text{m}$, $N_1 = 3.0 \times 10^{10}$,
 $\epsilon = 10 \text{ mm}\cdot\text{mrad}$
 Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$, $\sigma_z = 10.0 \mu\text{m}$, $N_2 = 1.0 \times 10^{10}$, $\epsilon = 0.093 \text{ mm}\cdot\text{mrad}$
 Distance between two beams : $115 \mu\text{m}$;
 Plasma Density : $1.0 \times 10^{17} \text{ cm}^{-3}$ (Field Ionized Li Plasma)



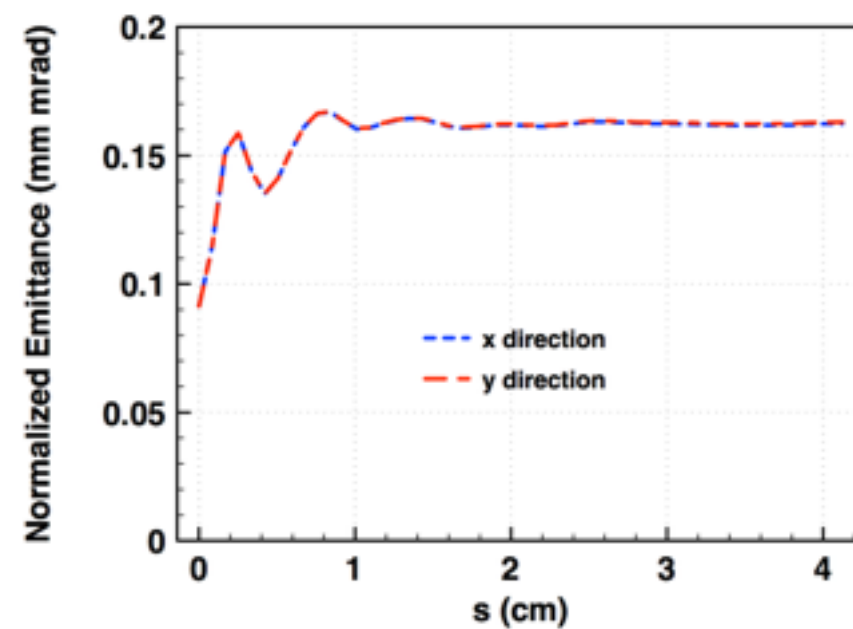
* W. An et al. Phys. Rev. Lett. 118, 244801, 2017 $\Delta E_z = \int dr \partial F_f / \partial \xi \approx \Delta r \Delta F_f / 2 \Delta \xi$

The Evolution of the electron beam's emittance

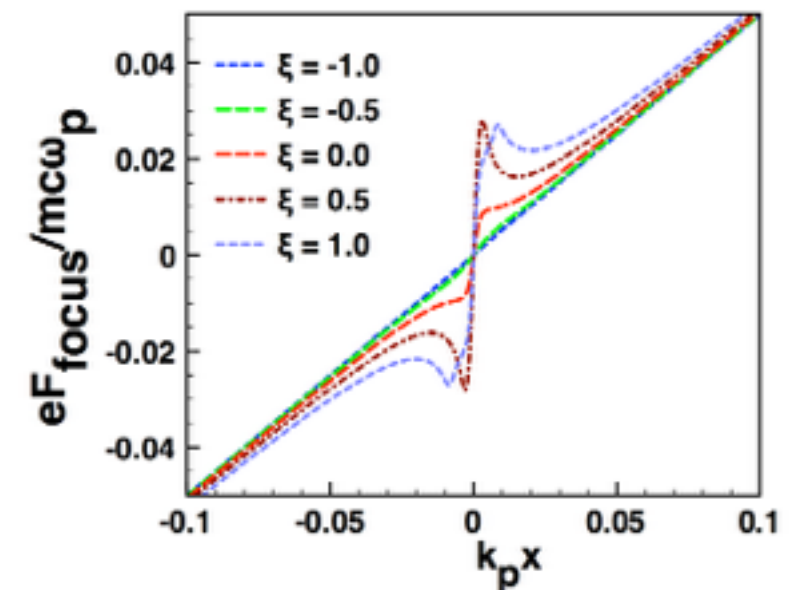
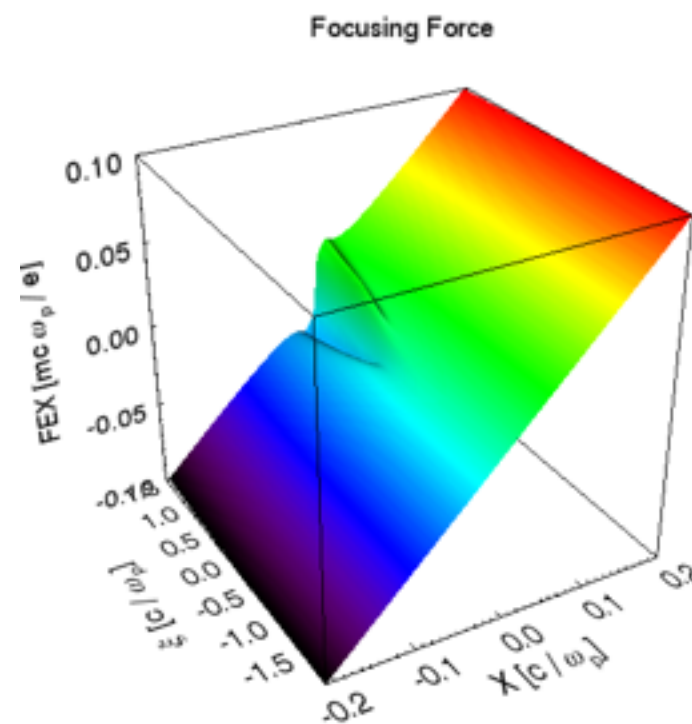
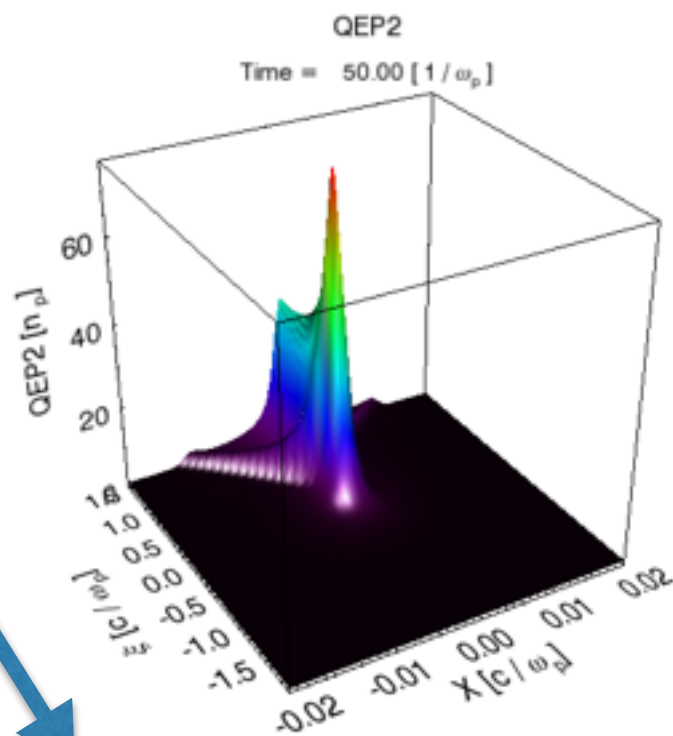
Li⁺



H⁺



Li⁺

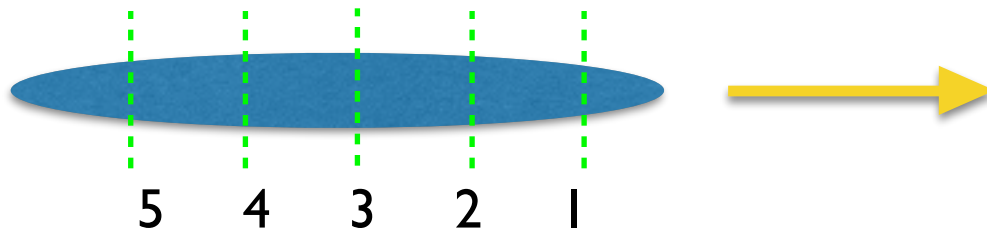


Emittance growth of the trailing beam

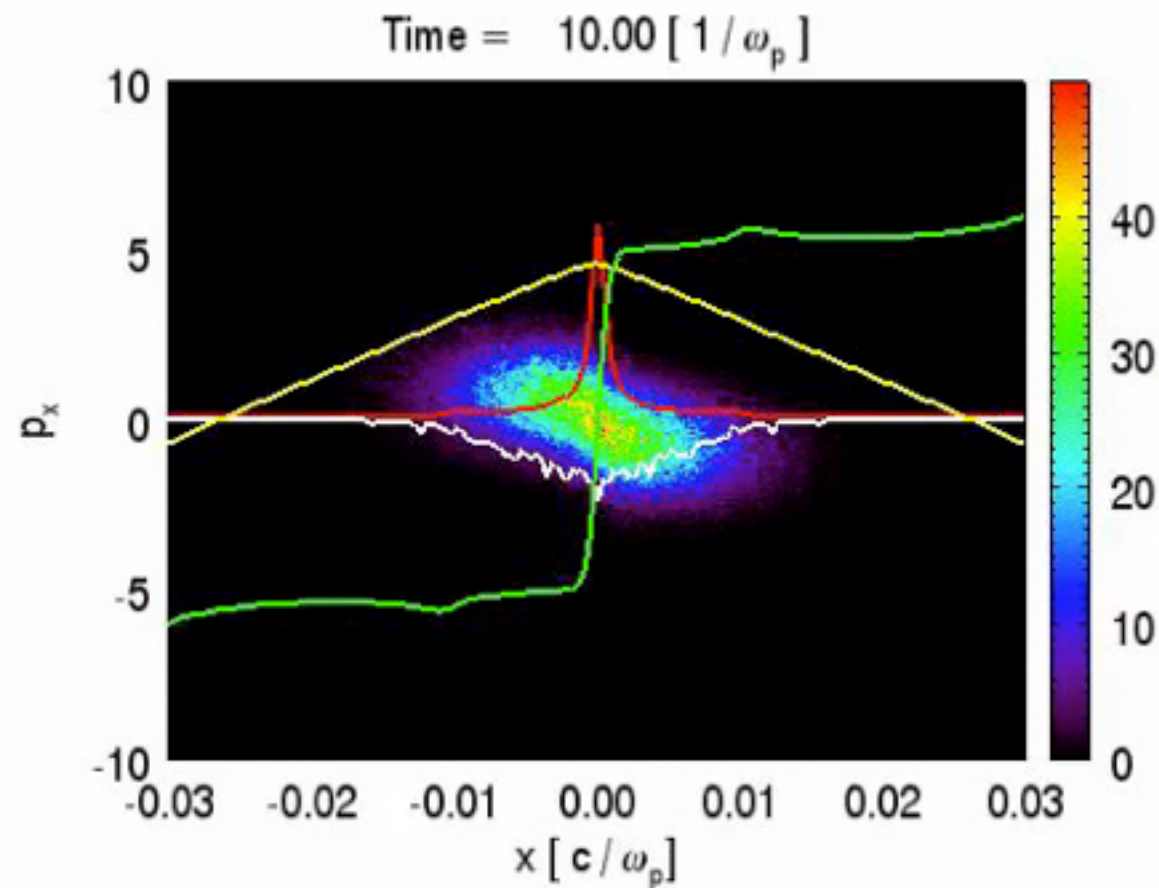
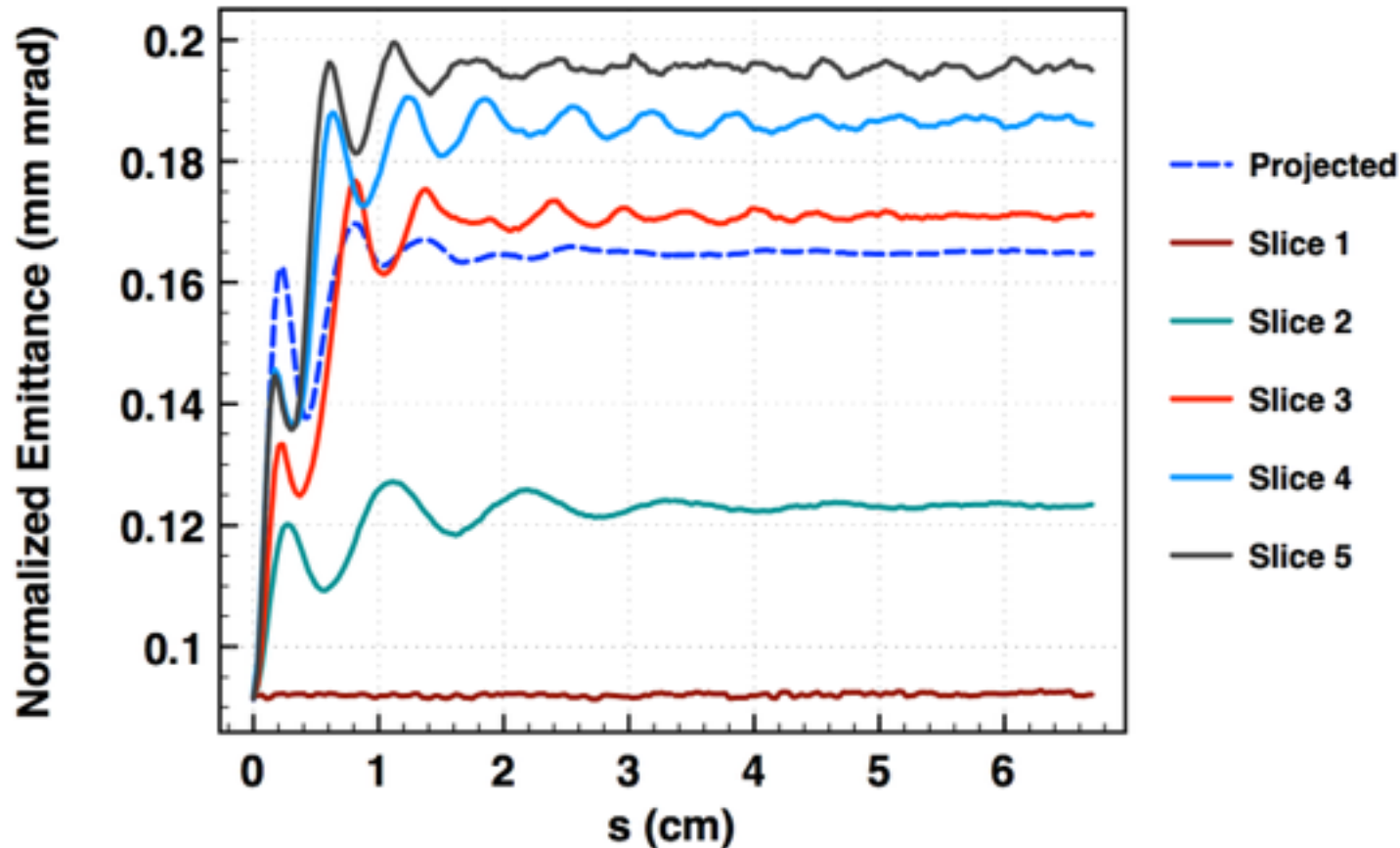
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Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$ ($0.006 k_p^{-1}$) , $\sigma_z = 10.0 \mu\text{m}$, $N_2 = 1.0 \times 10^{10}$, $\epsilon = 0.093 \text{ mm}\cdot\text{mrad}$

Distance between two beams : $115 \mu\text{m}$ Plasma Density : $1.0 \times 10^{17} \text{ cm}^{-3}$ (H Plasma)

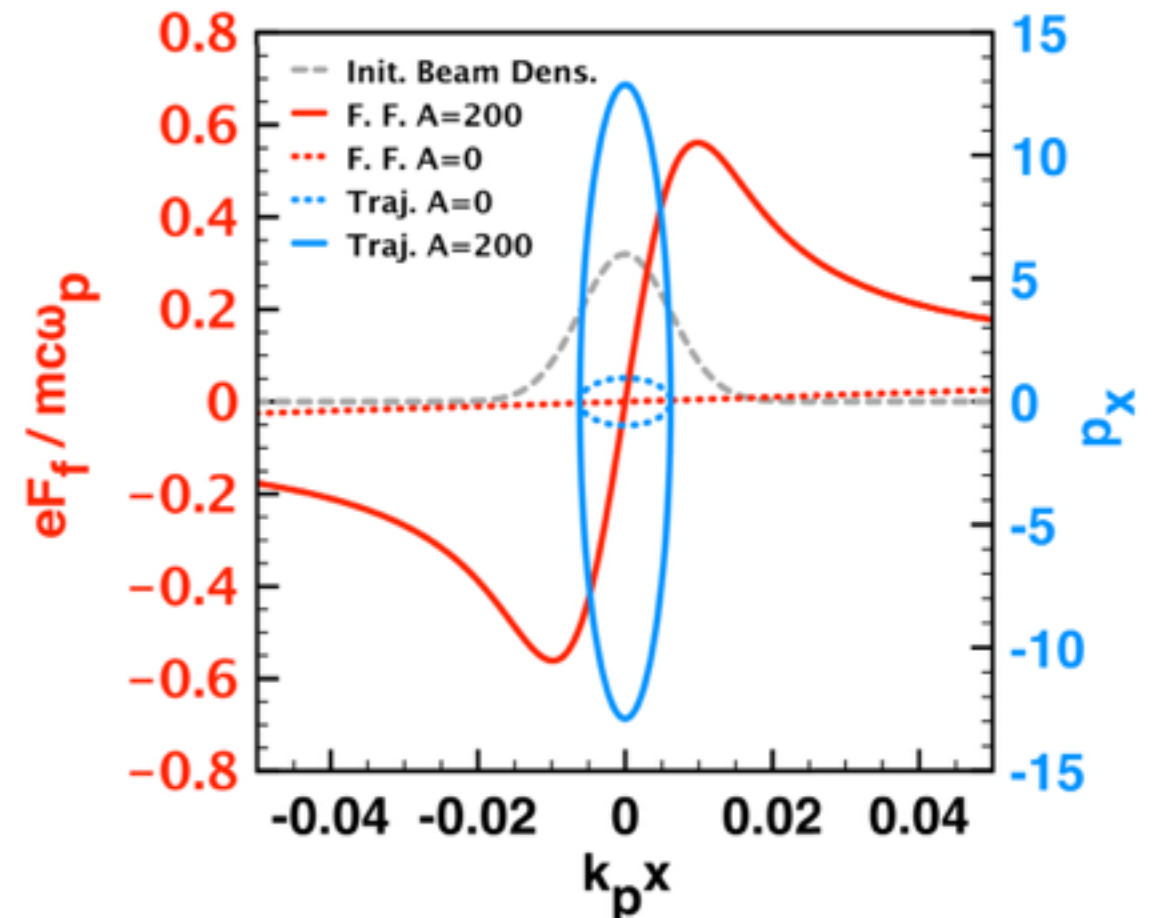
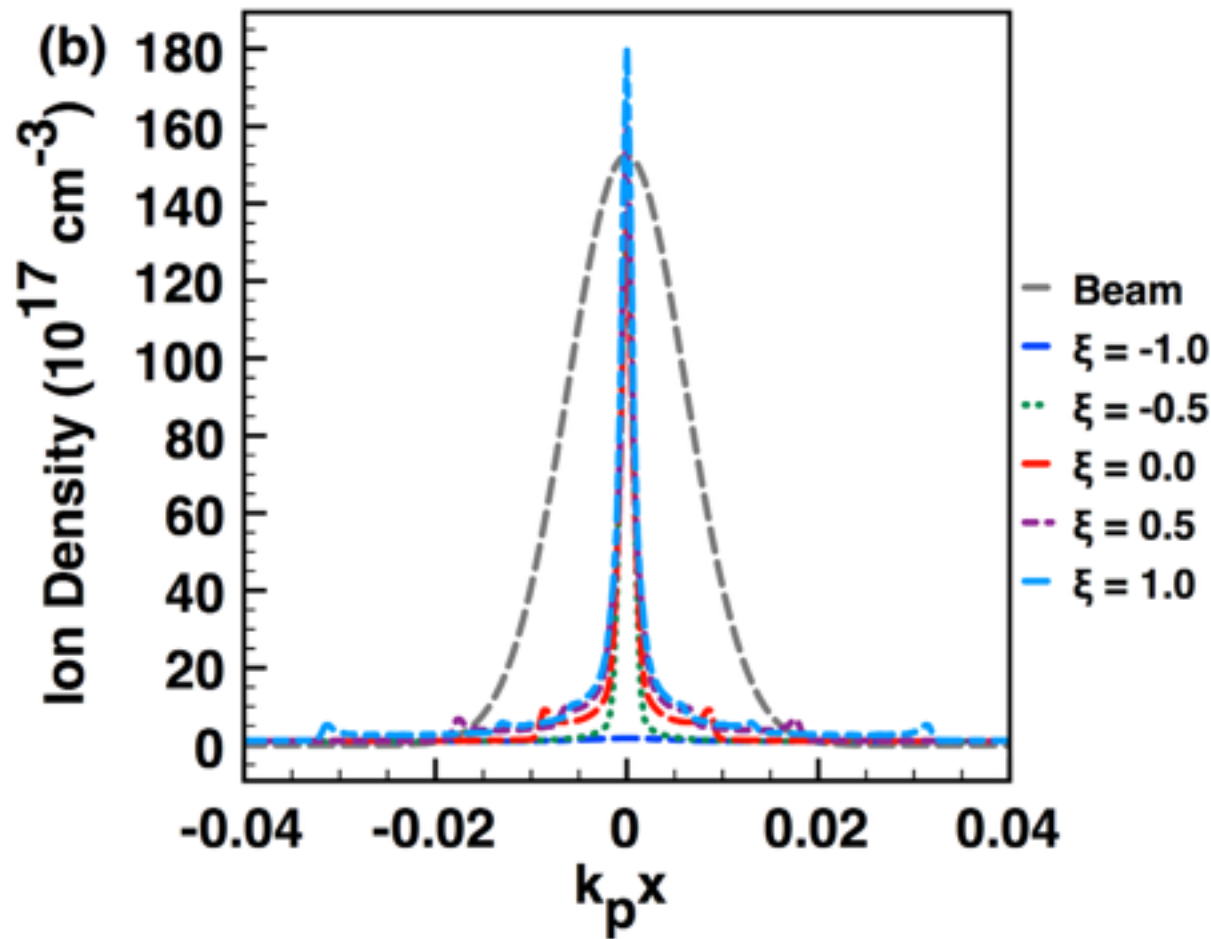


Slice 5

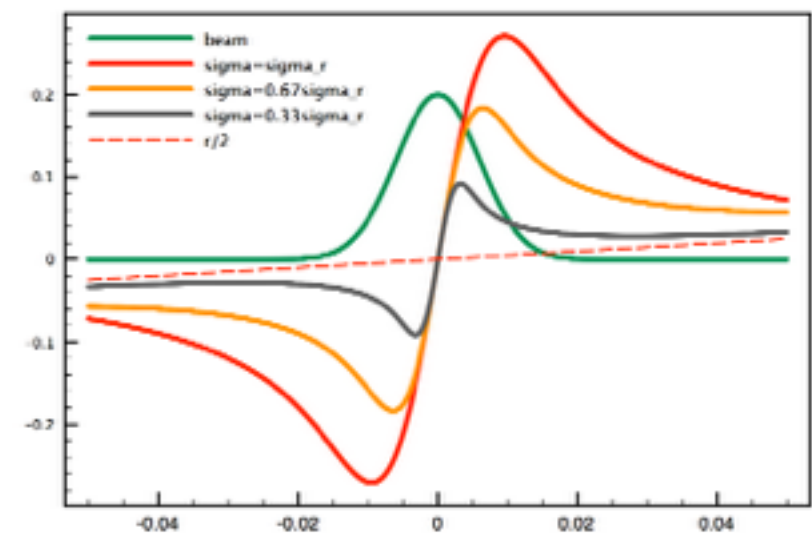


Ion density peak

The Ion profile: $1 + A \exp(-r^2/2\sigma^2)$



$$\epsilon_{Nf}/\epsilon_{N0} \propto p_{\max} \propto A^{0.5}$$

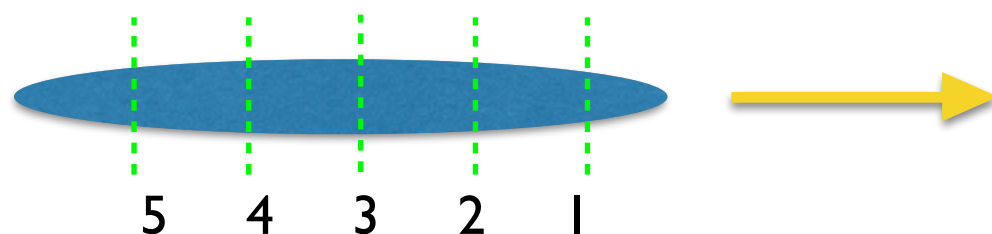


Emittance growth of the trailing beam

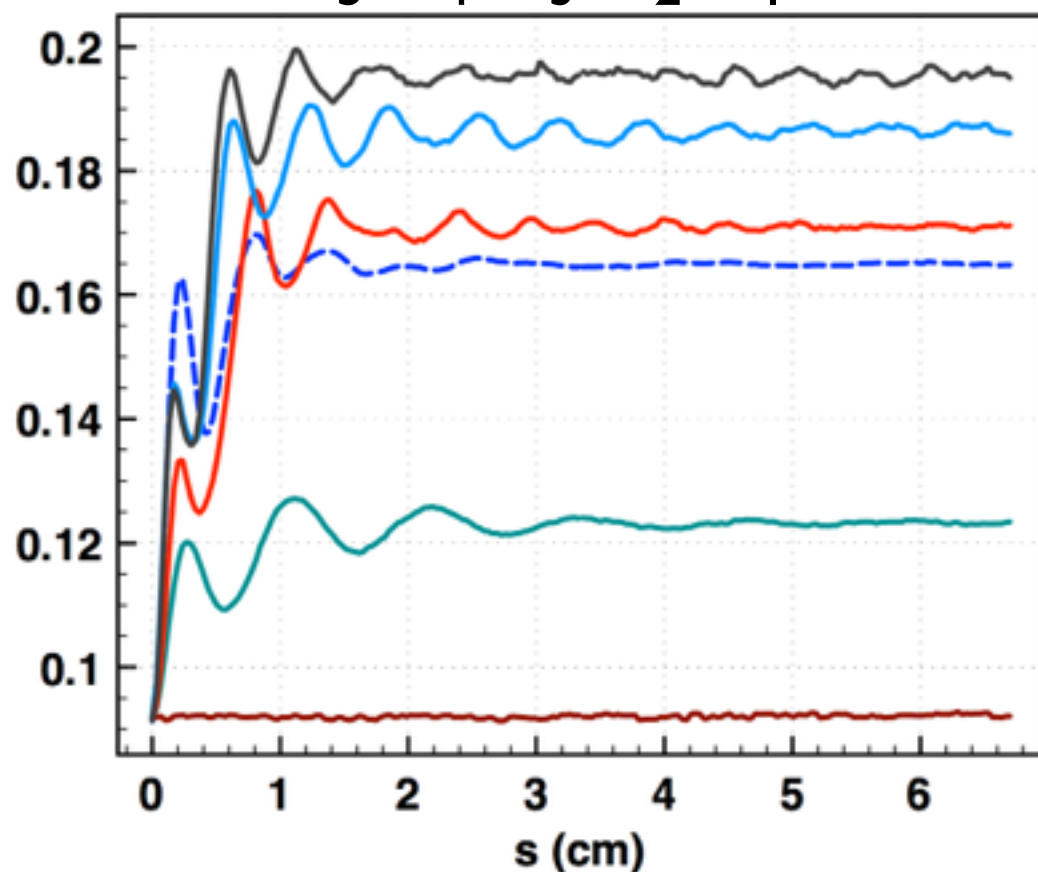
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Trailing Beam: $\sigma_r = 0.1 \mu\text{m}$ ($0.006 k_p^{-1}$) , $\sigma_z = 10.0 \mu\text{m}$, $N_2 = 1.0 \times 10^{10}$, $\epsilon = 0.093 \text{ mm}\cdot\text{mrad}$

Distance between two beams : $115 \mu\text{m}$ Plasma Density : $1.0 \times 10^{17} \text{ cm}^{-3}$ (H Plasma)



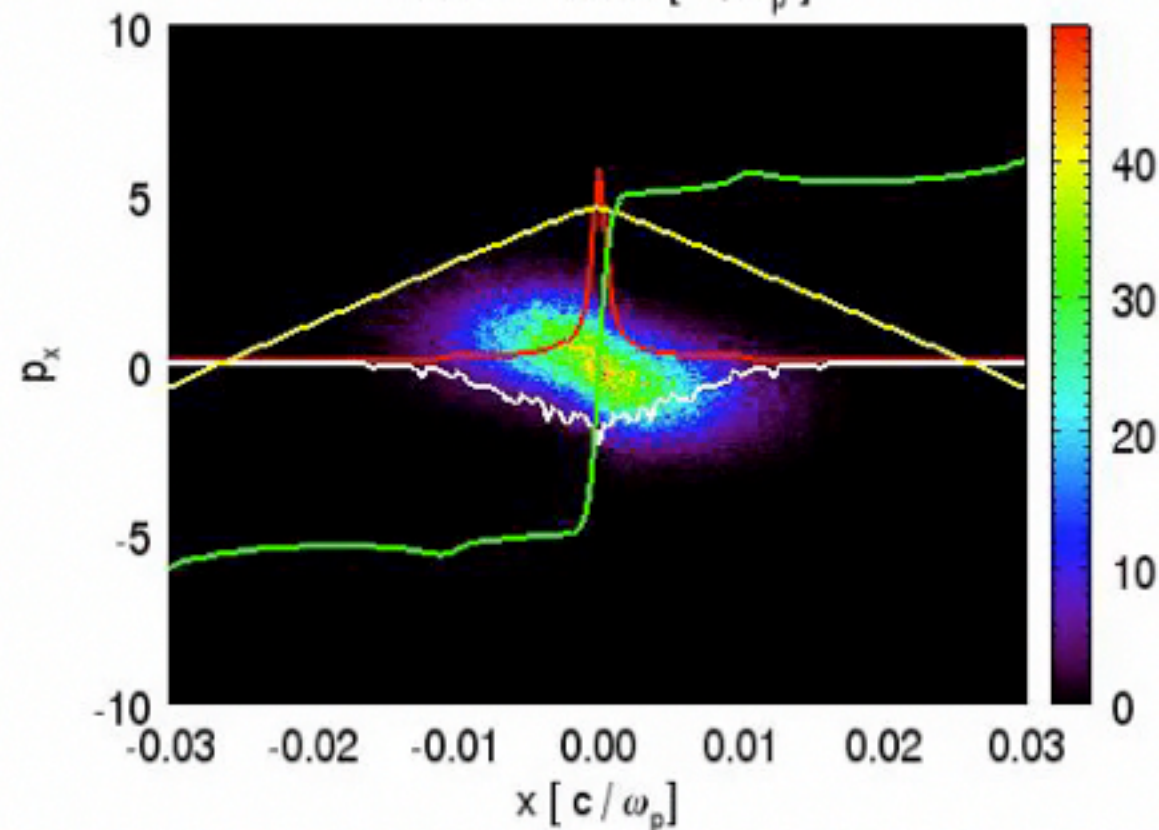
Normalized Emittance (mm mrad)



- Projected
- Slice 1
- Slice 2
- Slice 3
- Slice 4
- Slice 5

Slice 5

Time = $10.00 [1/\omega_p]$

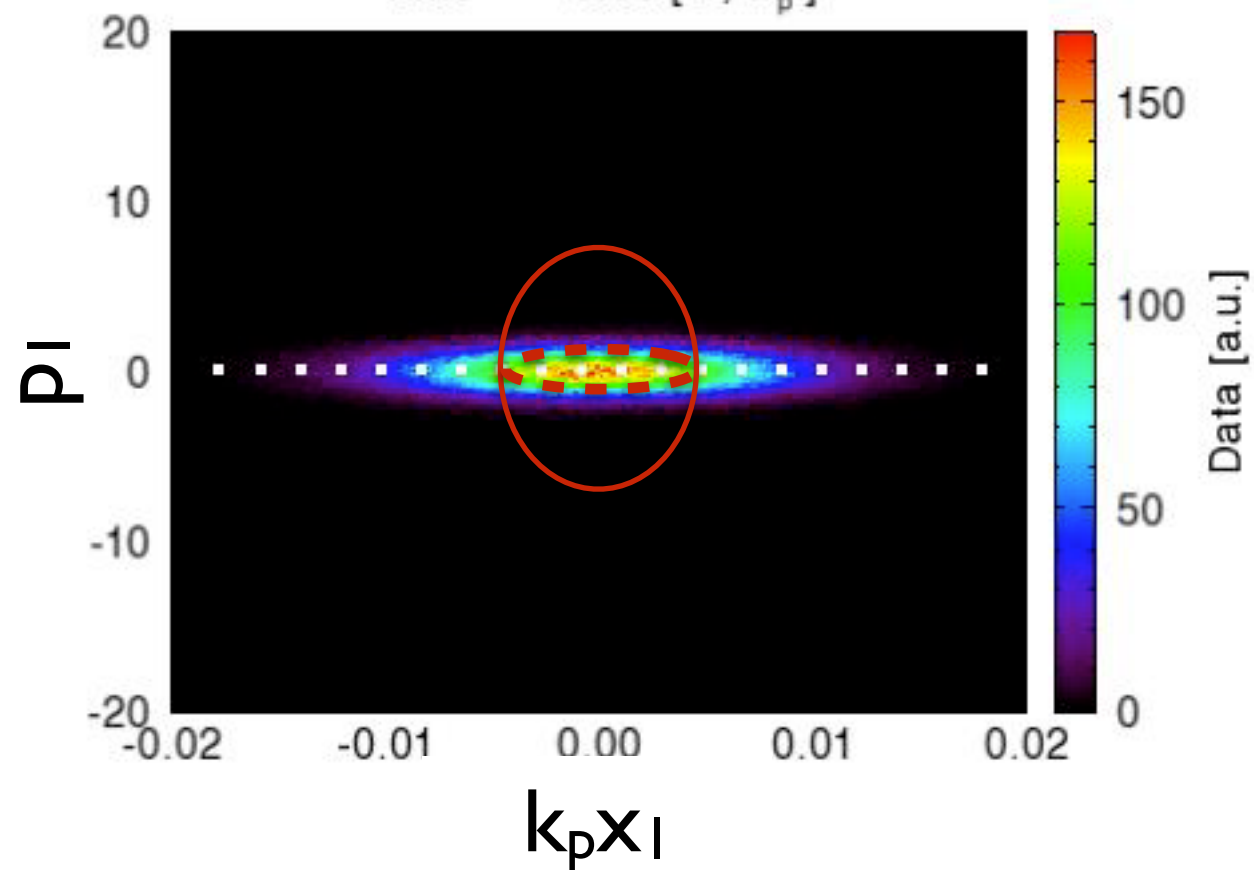


$$\rho_{\text{max}} / \rho_0 \sim 5 \gg \epsilon_{Nf} / \epsilon_{N0} \sim 2$$

Particle's trajectory

Phase Space of $p_{\perp} \times I$

Time = 0.00 [1 / ω_p]



ρ_{\max}

$\rho_{\text{time_averaged}}$

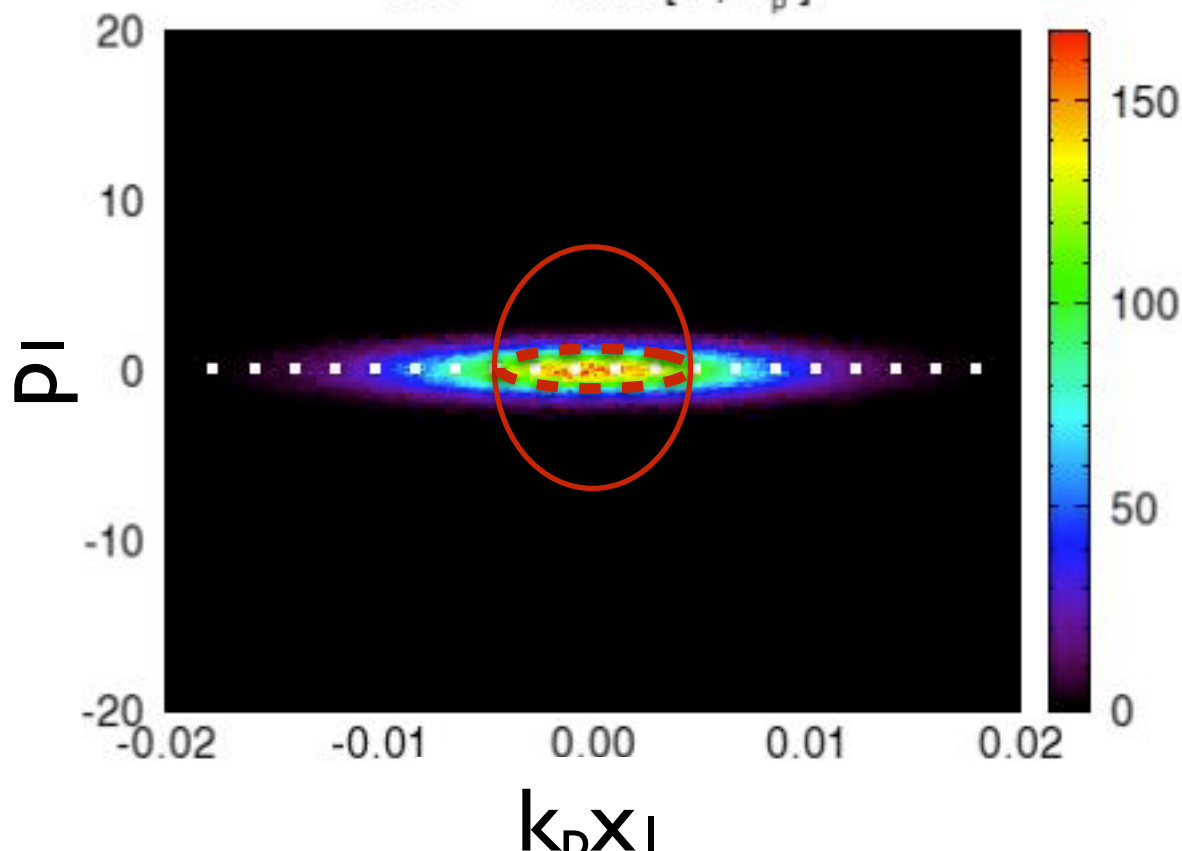
Emittance in the equilibrium state

$$\epsilon_N = \sqrt{\langle x^2 \rangle \langle p_x^2 \rangle - \langle xp_x \rangle^2}$$

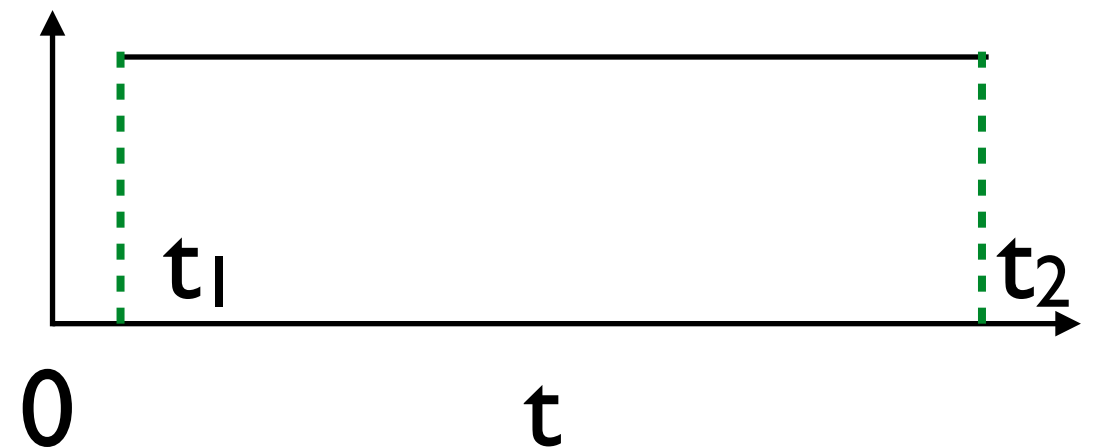
$$\epsilon_{Nf} = \sqrt{\langle x^2 \rangle_f \langle p_x^2 \rangle_f}$$

Phase Space of $p|x|$

Time = 0.00 [1/ω_p]



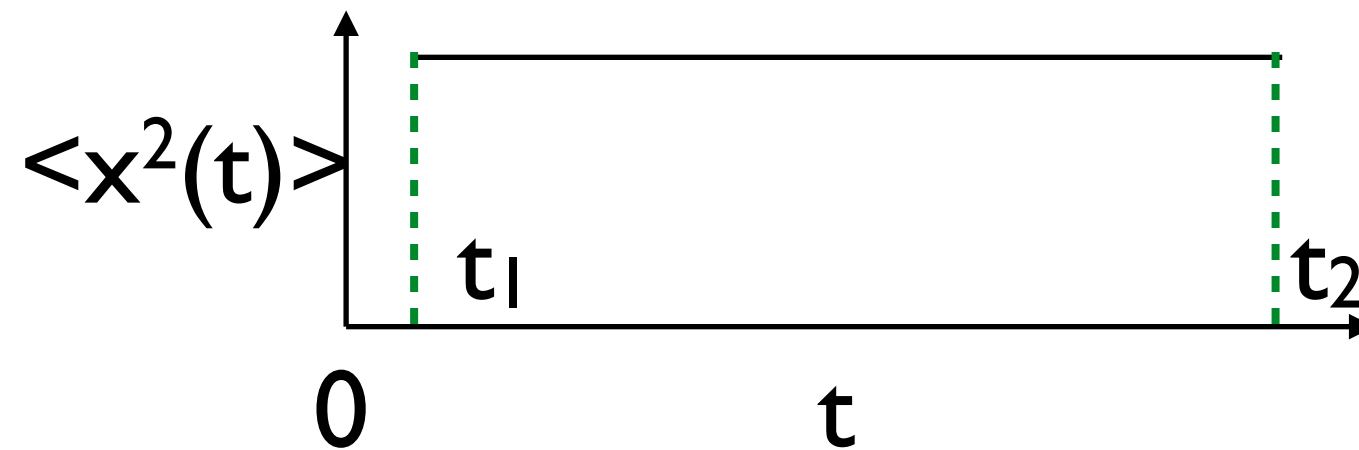
$\langle x^2 \rangle$



Data [a.u.]

$$\langle x^2 \rangle_f = \frac{1}{t_2 - t_1} \int_{t_1}^{t_2} \langle x^2 \rangle dt = \left\langle \frac{\int_{t_1}^{t_2} x^2 dt}{t_2 - t_1} \right\rangle$$

Emittance in the equilibrium state



$$\langle x^2 \rangle_f = \frac{1}{t_2 - t_1} \int_{t_1}^{t_2} \langle x^2 \rangle dt = \left\langle \frac{\int_{t_1}^{t_2} x^2 dt}{t_2 - t_1} \right\rangle$$

Each
particle

$$\frac{\int_{t_1}^{t_2} x^2 dt}{t_2 - t_1} = \frac{\int_{osc} x^2 dt}{T_{osc}} \quad \text{when } t_2 - t_1 \gg T_{osc}$$

$$\langle x^2 \rangle_f = \left\langle \frac{\int_{osc} x^2 dt}{T_{osc}} \right\rangle$$

$$\langle x^2 \rangle_f = \left\langle \frac{\int_{osc} x^2 dt}{T_{osc}} \right\rangle$$

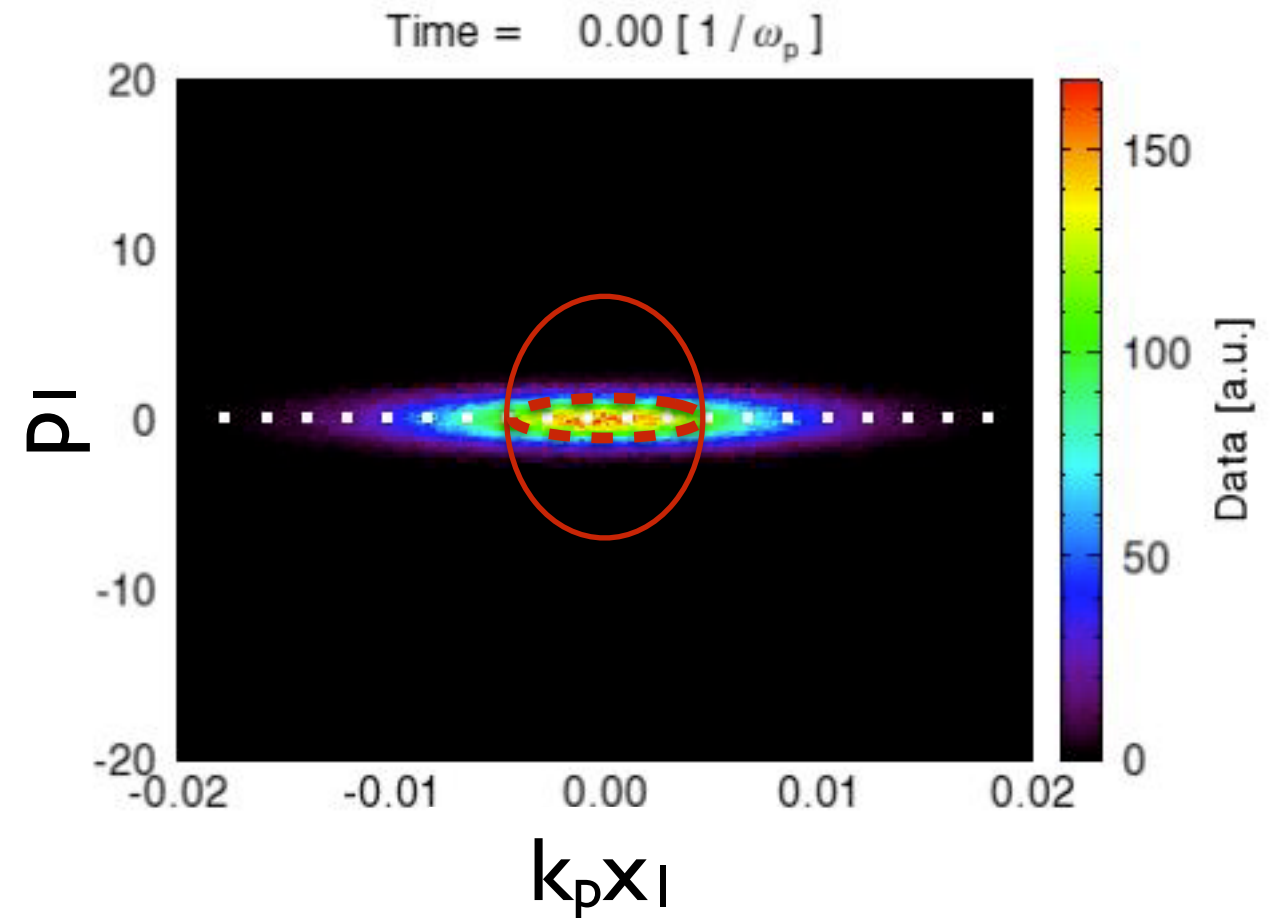
Sort the particle according to x_0

$$\langle x^2 \rangle_f = \left\langle \frac{\int_{x_0} x^2 dt}{T_{x_0}} \right\rangle$$

$$\langle x^2 \rangle_f = \left\langle \frac{\int_{x_0} x^2 dt}{T_{x_0}} \right\rangle = \frac{1}{N} \int_0^\infty N_{x_0} \frac{\int_{x_0} x^2 dt}{T_{x_0}} dx_0$$

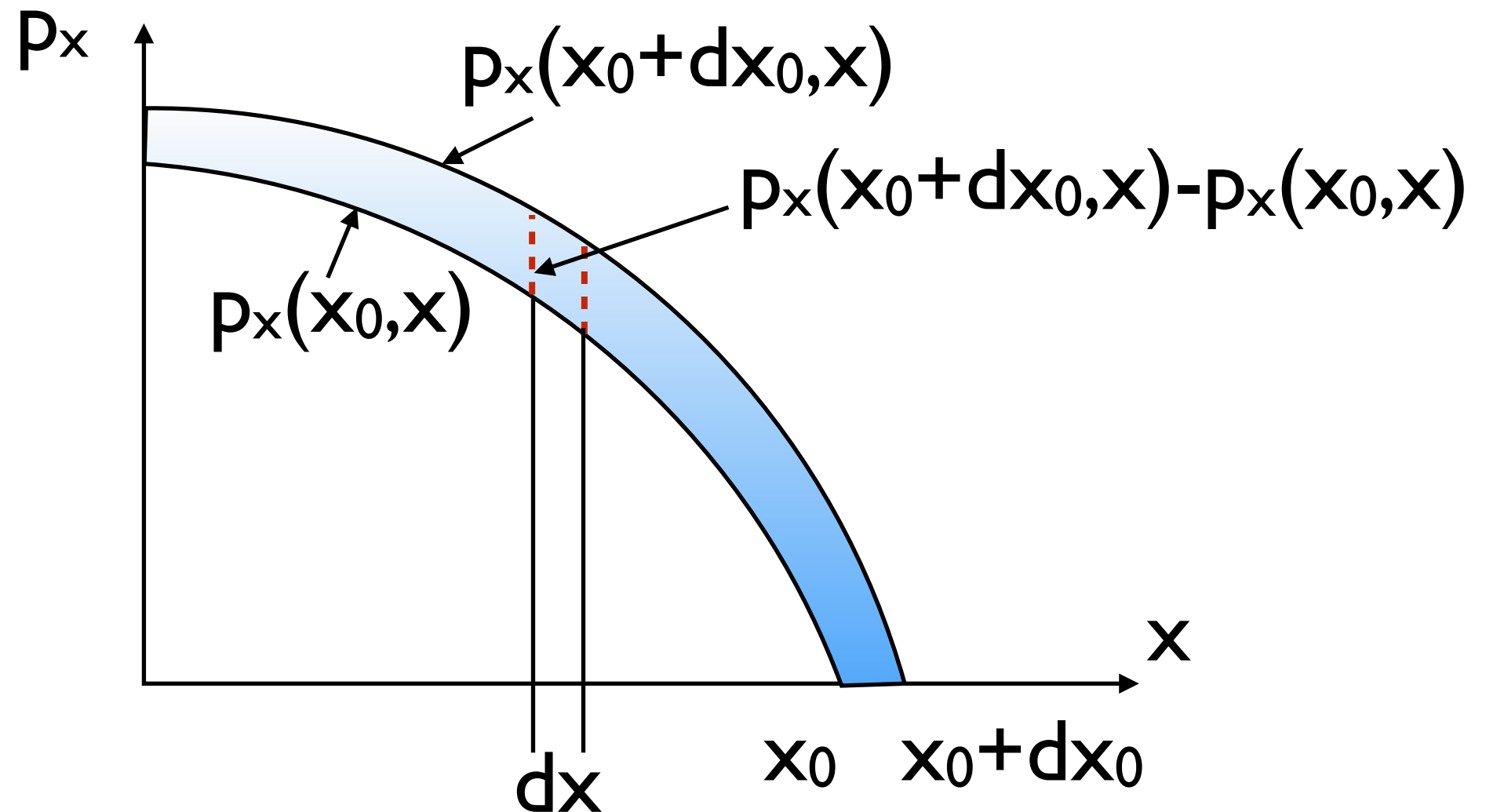
N_{x_0} does not change because phase space trajectories do not cross with each other.

Phase Space of $p|x|$



Calculating N_{x_0}

$$N_{x_0} = 4 \int_0^{x_0} f_0(x, p_x) \frac{dp_x}{dx_0} dx$$



$$\langle x^2 \rangle_f = \langle X_{\text{ave}}^2 \rangle = \frac{1}{N} \int_0^\infty dx_0 N_{x_0} X_{\text{ave}}^2(x_0)$$

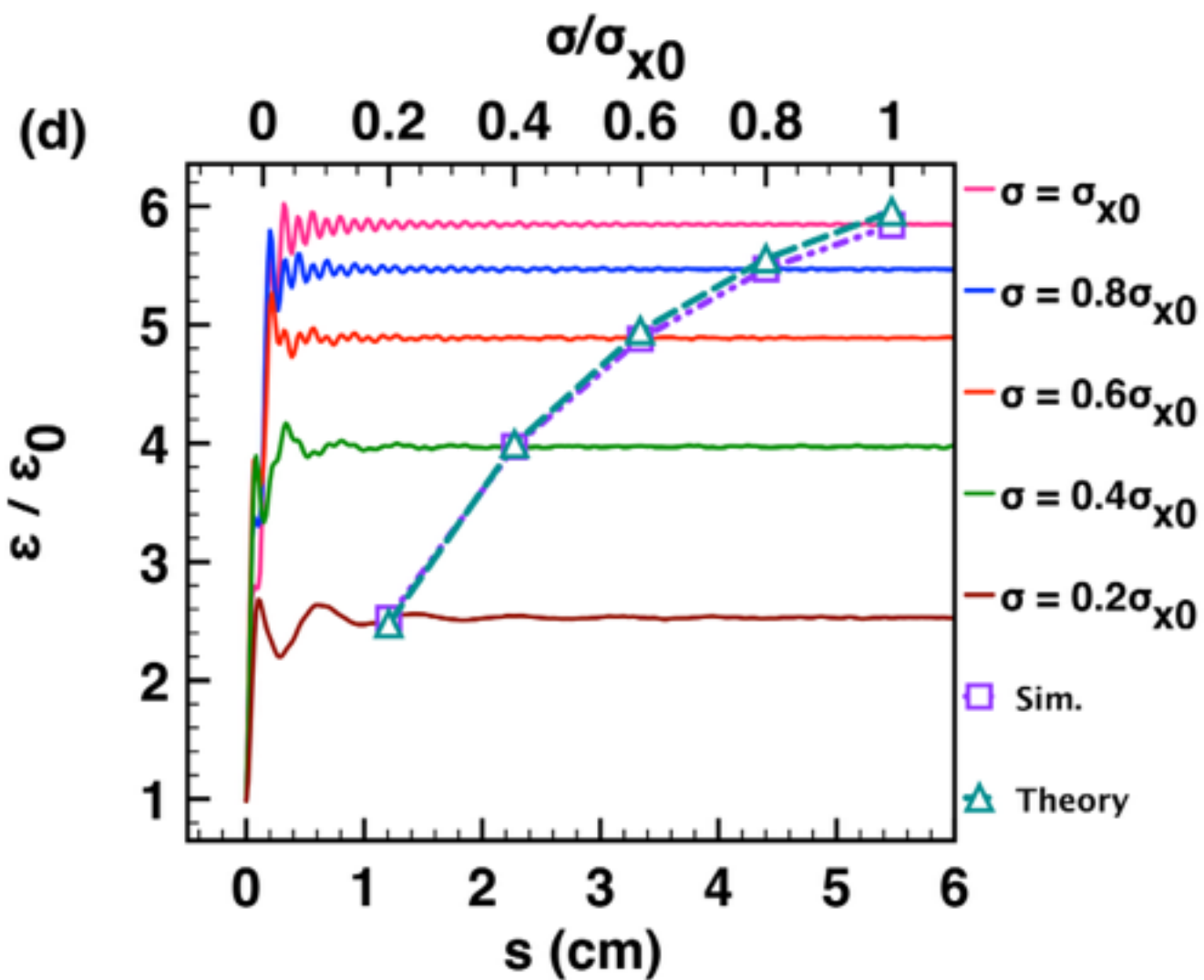
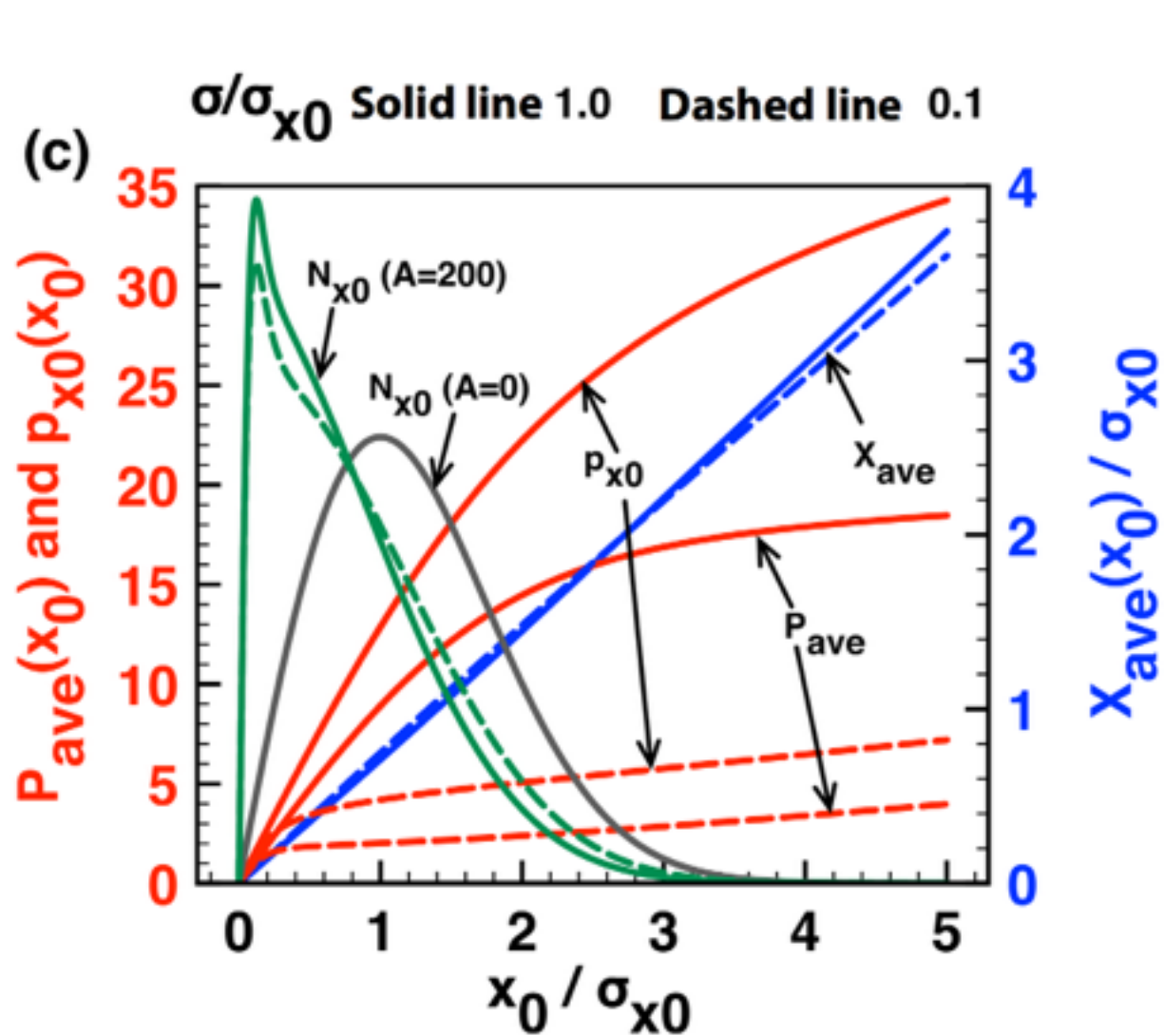
$$\langle p_x^2 \rangle_f = \frac{1}{N} \int_0^\infty dx_0 N_{x_0} P_{\text{ave}}^2(x_0)$$

$$N_{x_0} = 4 \int_0^{x_0} f_0(x, p_x) \frac{dp_x}{dx_0} dx$$

$$X_{\text{ave}}^2 = \frac{\int_0^{x_0} \frac{dx}{v_x} x^2}{\int_0^{x_0} \frac{dx}{v_x}} = \frac{\int_0^{x_0} dx x^2 / \sqrt{(\psi(x, \xi) - \psi(x_0, \xi))}}{\int_0^{x_0} dx / \sqrt{(\psi(x, \xi) - \psi(x_0, \xi))}}$$

$$P_{\text{ave}}^2 = \gamma \frac{\int_0^{x_0} dx \sqrt{(\psi(x, \xi) - \psi(x_0, \xi))}}{\int_0^{x_0} dx / \sqrt{(\psi(x, \xi) - \psi(x_0, \xi))}}$$

$$\partial p_x / \partial x_0 = \sqrt{\gamma} F_f(x_0) / \sqrt{2(\psi(x) - \psi(x_0))}$$

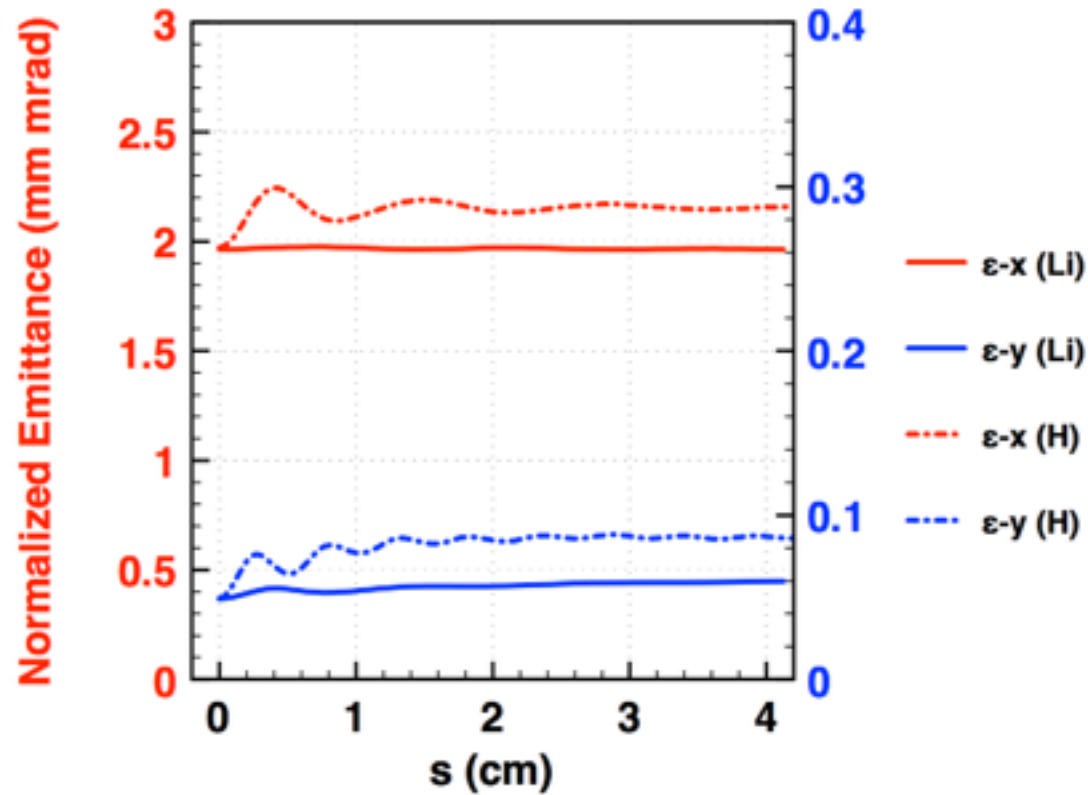


Trailing Beam: $\sigma_z = 10.0 \mu\text{m}$, $N = 1.0 \times 10^{10}$,

$$\sigma_x / \Delta_{\perp} = 75.9$$

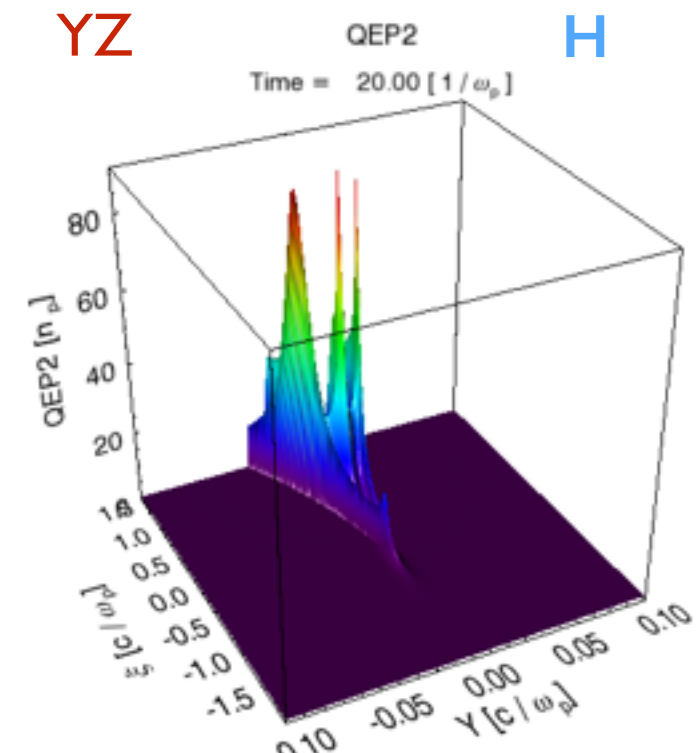
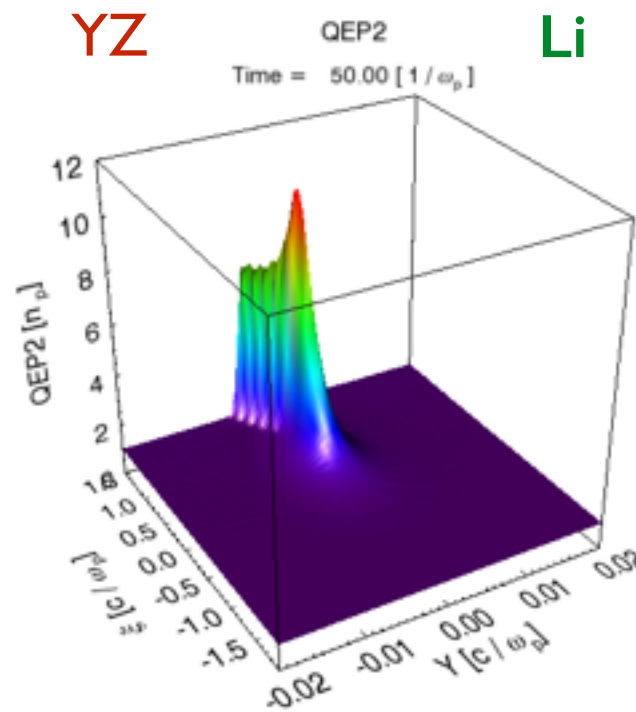
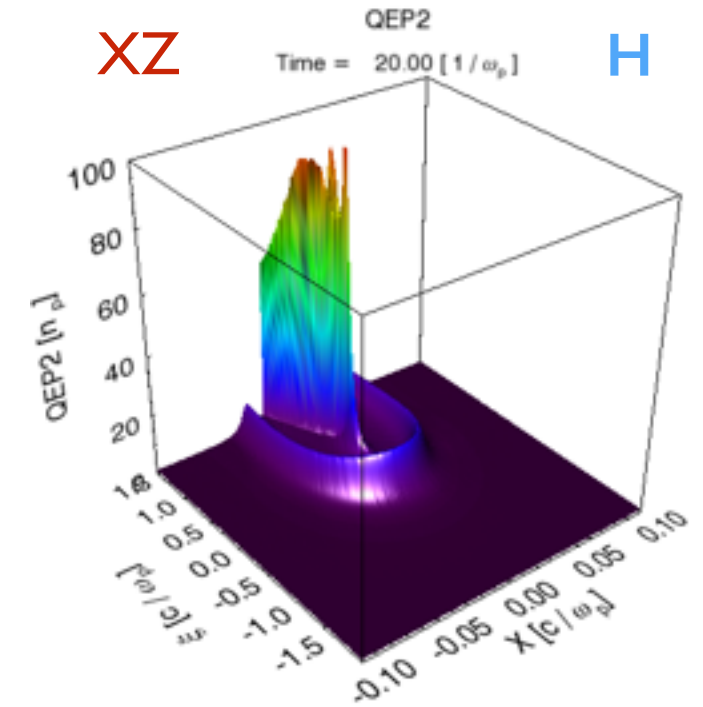
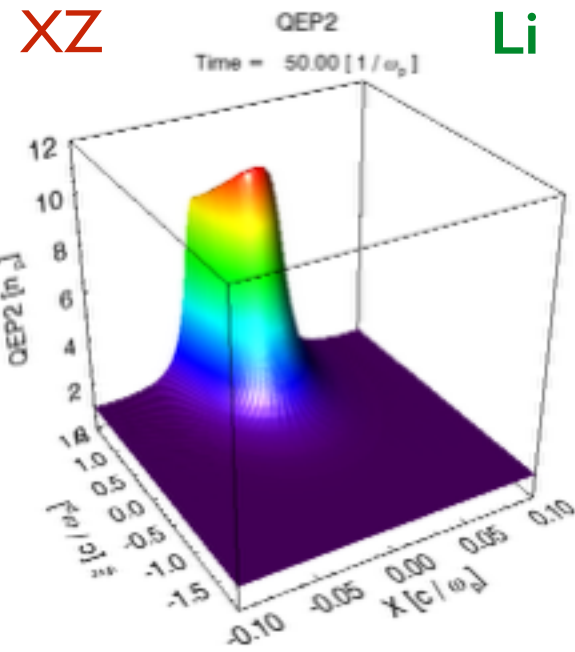
$$\sigma_y / \Delta_{\perp} = 12.0$$

$\sigma_x = 0.463 \mu\text{m}$, $\epsilon_{Nx} = \mathbf{2.0 \text{ mm}\cdot\text{mrad}}$, $\sigma_y = 0.0733 \mu\text{m}$, $\epsilon_{Ny} = \mathbf{0.05 \text{ mm}\cdot\text{mrad}}$
 $\mathbf{Y = 48923.7 (25 \text{ GeV})}$, Plasma Density : $1.0 \times 10^{17} \text{ cm}^{-3}$



In Li, the emittance in x does not change, and in y direction it only increase by 20%.

In H, the emittance in x increase by 10%, and in y direction it increases by 70%.



The emittance growth can be mitigated.

